

Microscopic structure of interfaces in condensed matter

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*Chair of Quantum Physics and
Nanoelectronics*

Nanotechnology: flourishing diversity

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Nanotechnologies in MIET:

- **Growth technologies: molecular-beam epitaxy, arc plasma coating, electro-chemical nanostructure formation, porous alumina anodic oxide, Ge nanocluster in Si matrix formation using ion implantation.**
- **Scanning probe technologies: local oxidation based nanolithography, material placing, nanolithography cantilever creation.**

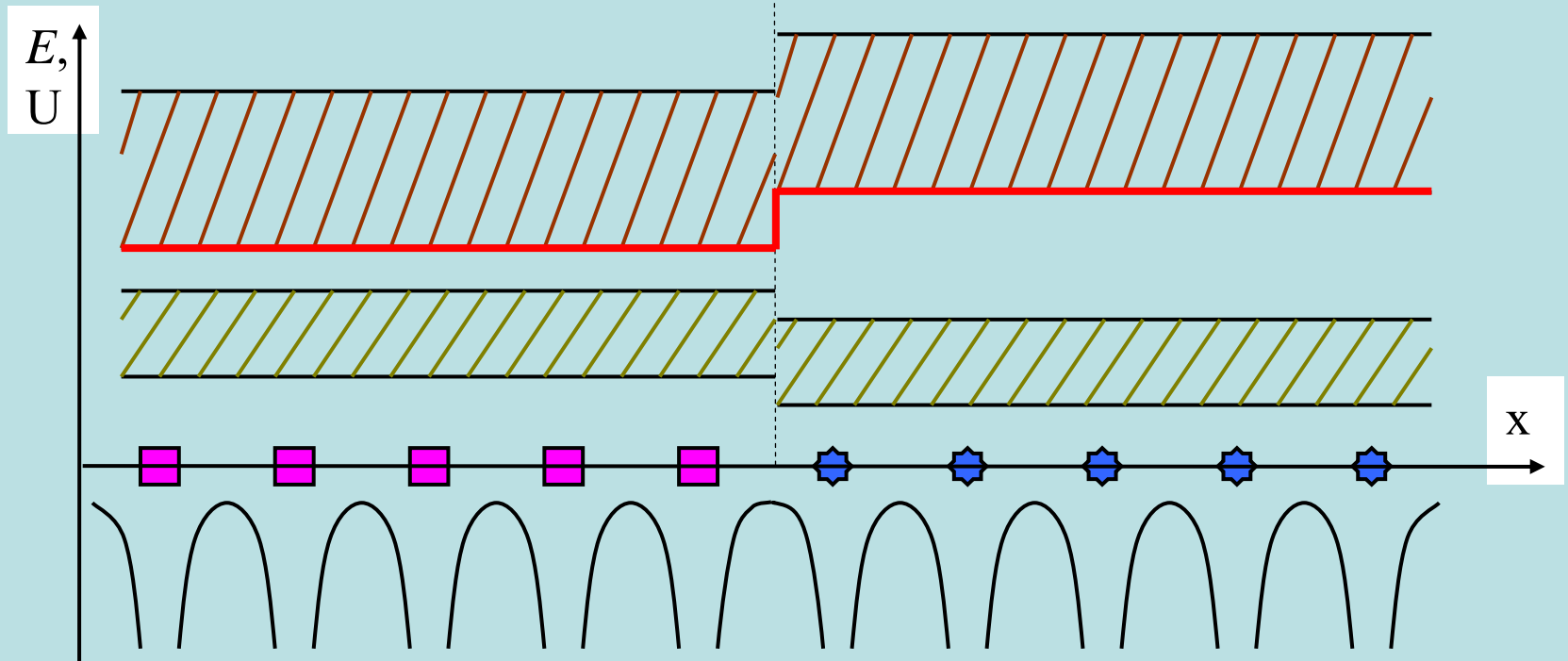
Nanomaterials and nanostructures for electronics

- semiconductor heterostructures based on A₃B₅ compounds, including resonant-tunneling heterostructures for ultra-high speed electronics;
- solid nanocoatings for various applications, including coatings for nanolithography, conducting and magnetic coatings for scanning probe microscopy cantilevers;
- carbon nanotubes for studies of quantum transport, development and creation of electronic devices based on quantum conductors;
- Ge quantum dots in Si matrix which were first obtained using ion implantation technique;
- nanoporous alumina oxide for optoelectronic applications as well as for terabit semiconductor memory cell creation;
- quasi-one-dimensional conductor based nanovaristors displaying quantum properties at room temperature, obtained by local anodic and current induced oxidation of various metals (Ti, Ta, Nb, Al)

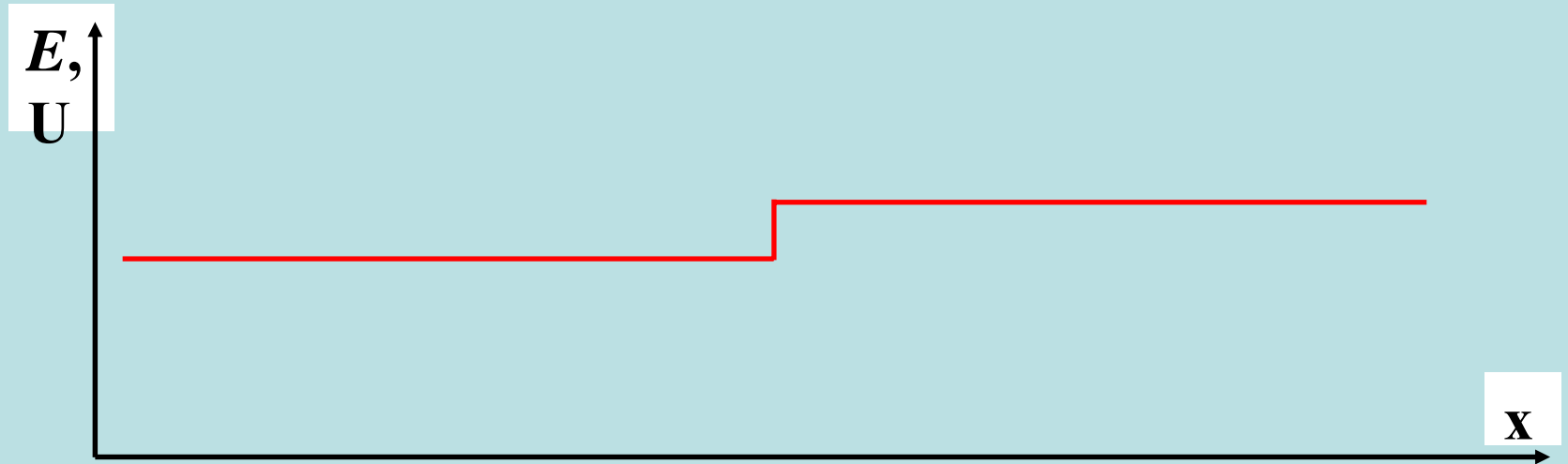
Nanomaterials and nanostructures for electronics

- semiconductor heterostructures based on A₃B₅ compounds, including resonant-tunneling heterostructures for ultra-high speed electronics;

Heterojunction

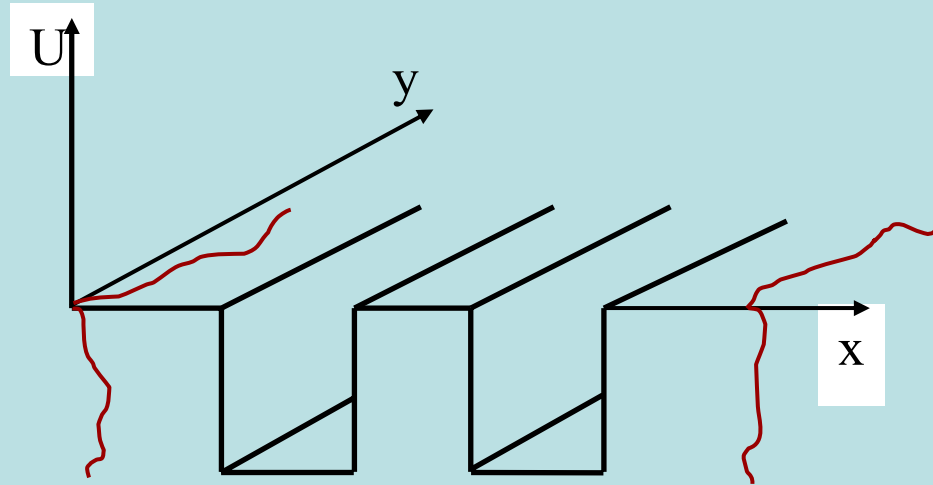


Envelope Function (or Effective Mass) Method

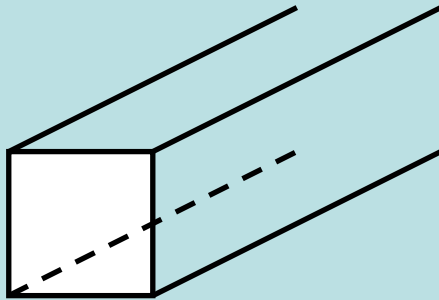


$$\left[-\frac{\hbar^2 \nabla^2}{2m^*} + U(\mathbf{r}) \right] \psi(\vec{r}) = E \psi(\vec{r})$$

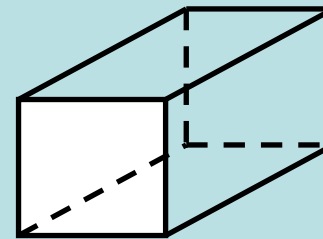
1. Quantum Wells and Superlattices



2. Quantum Wires



3. Quantum Dots

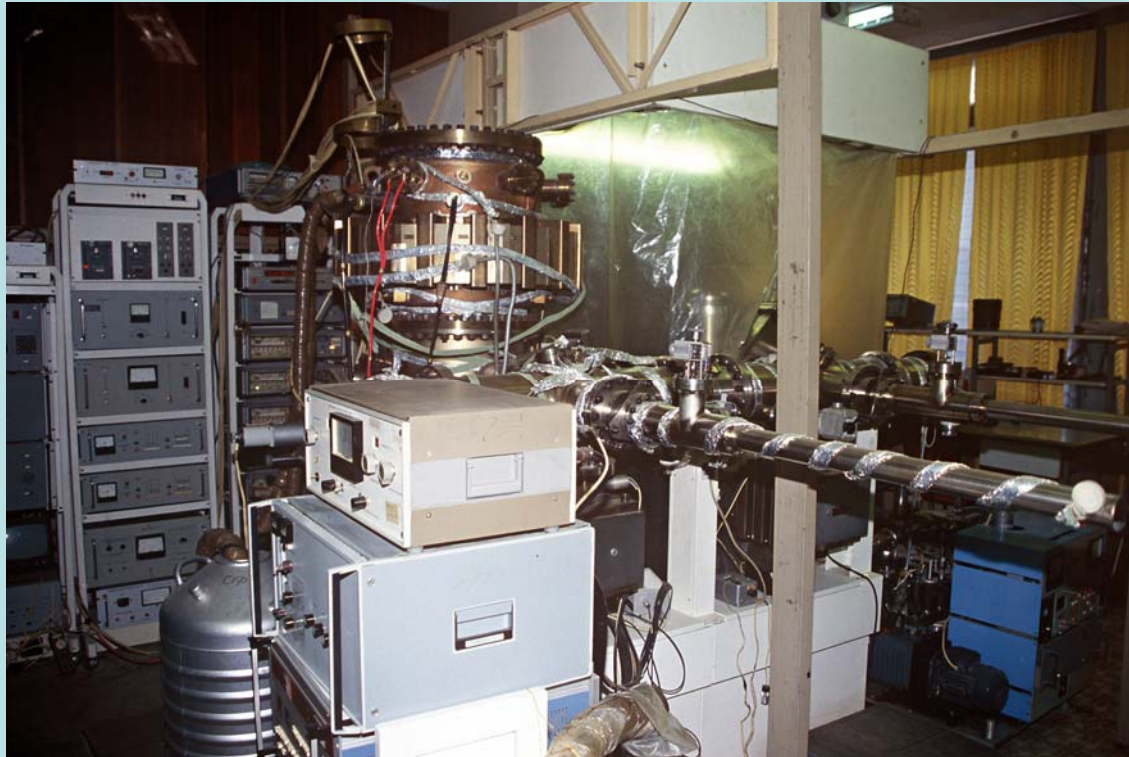


- Heterostructures for high-speed IC applications
- Nanoelectronic elements for ultra-high-speed ICs
- Heterostructures for fundamental researches (microscopic structure of heterojunctions)

- Heterostructures for high-speed IC applications
- Nanoelectronic elements for ultra-high-speed ICs
- Heterostructures for fundamental researches (“dark matter” in crystals)

Applications

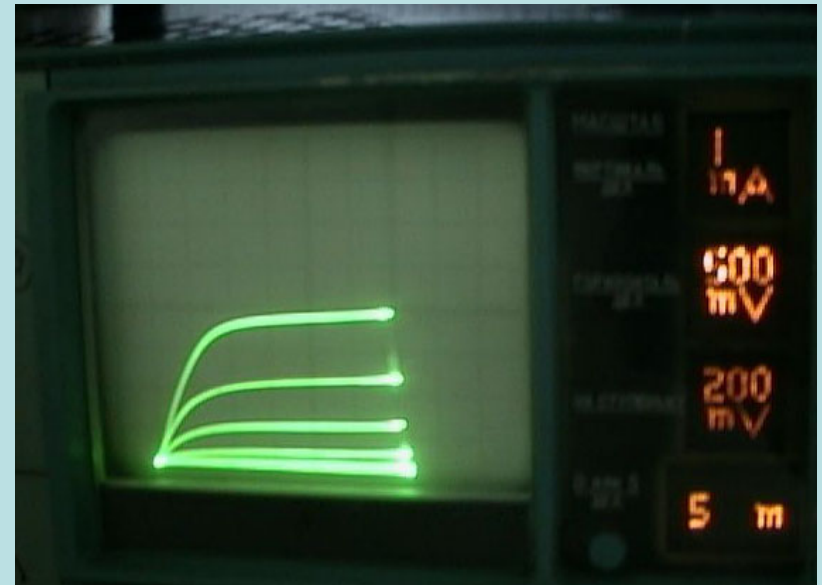
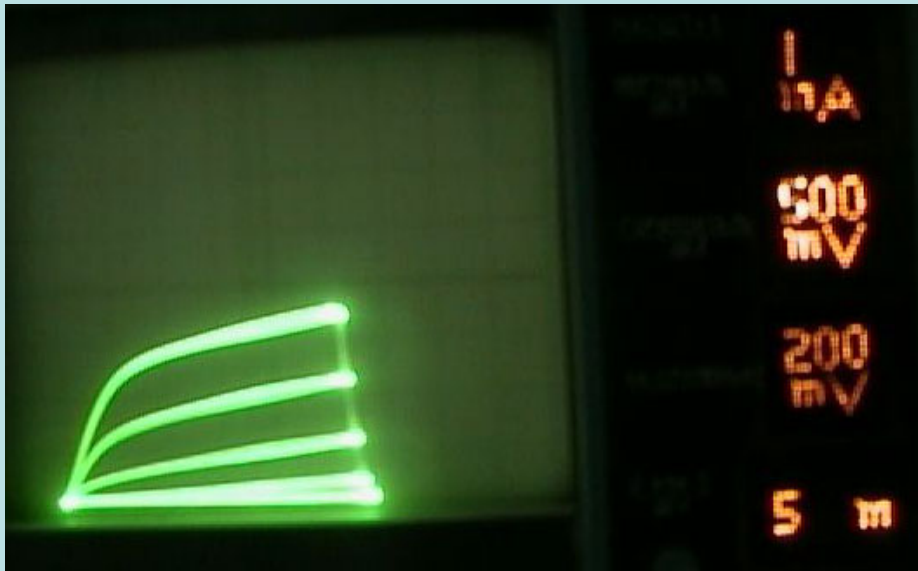
Molecular-beam epitaxy machine



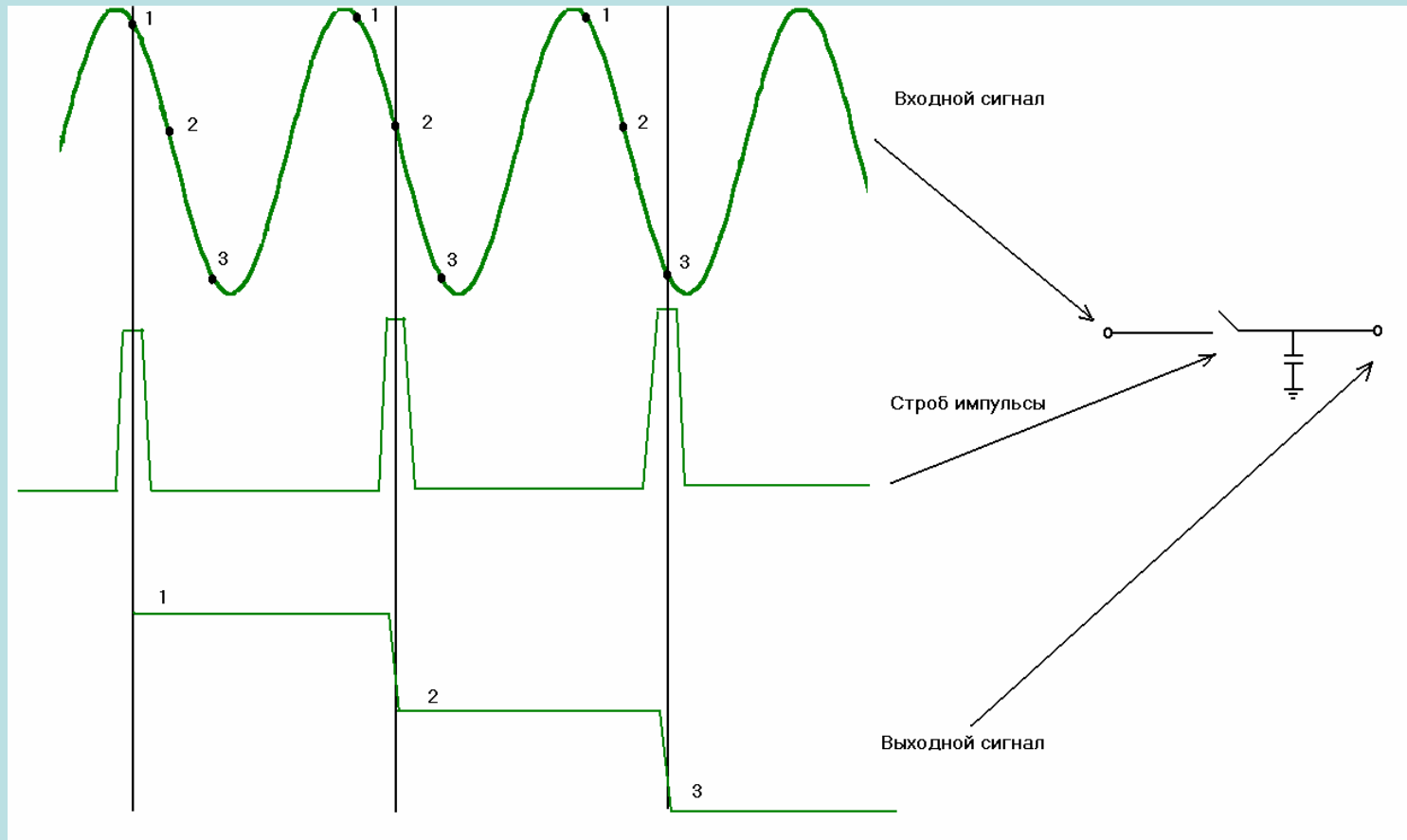
General view of A3B5 compounds technology clean-room facilities



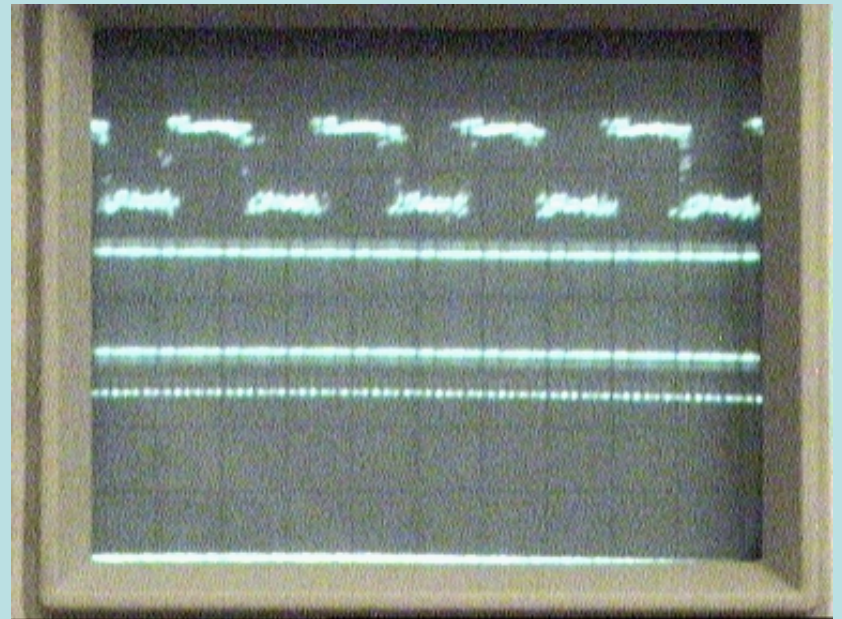
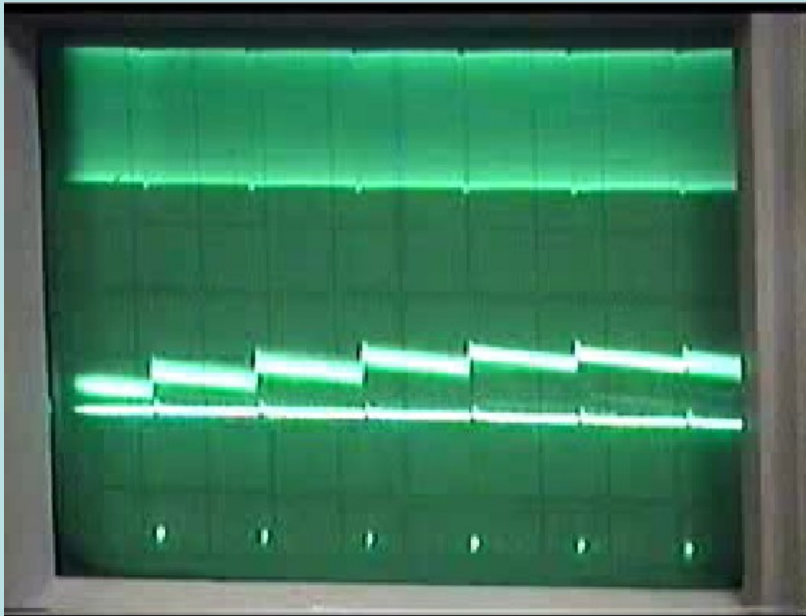
Comparison of FET based on doped epitaxial GaAs structure and heterostructure FET



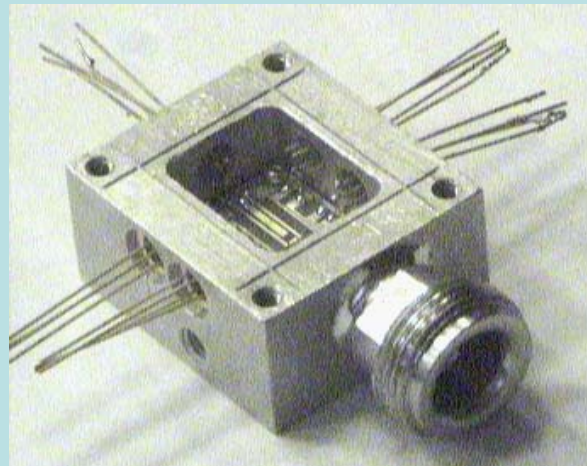
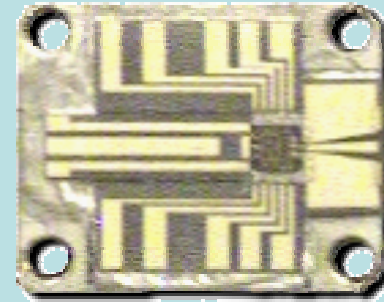
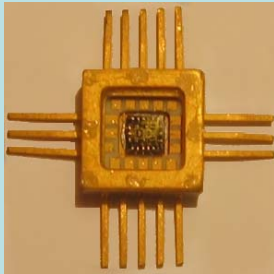
Basic principal of sampler performance



Input signal expansion mode



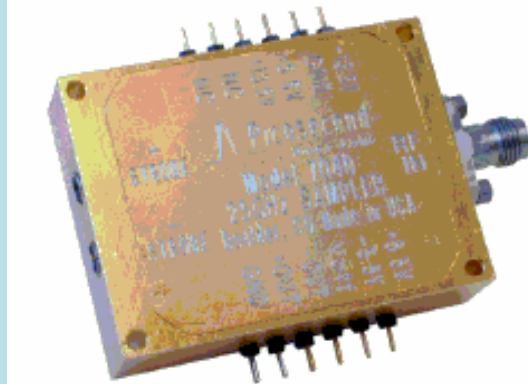
EXAMPLES OF RF ICs AND MODULES



Analogs by Picosecond Pulse Labs 2004-2005 г.

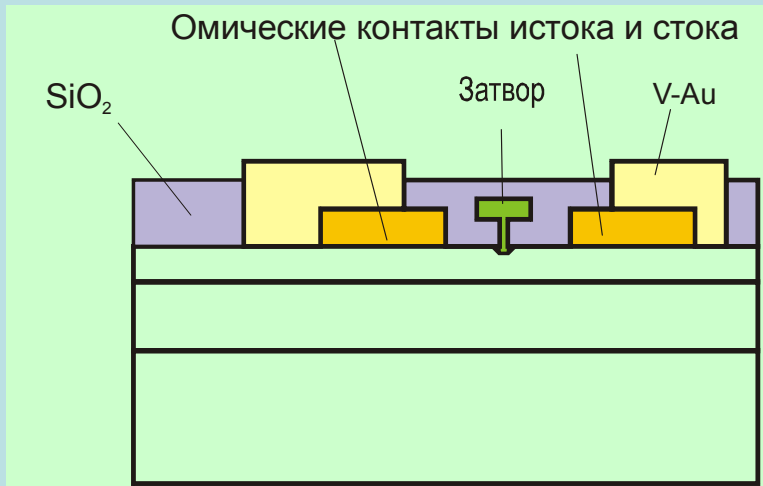


**MODEL 7040 25 GHz
SAMPLER MODULE**

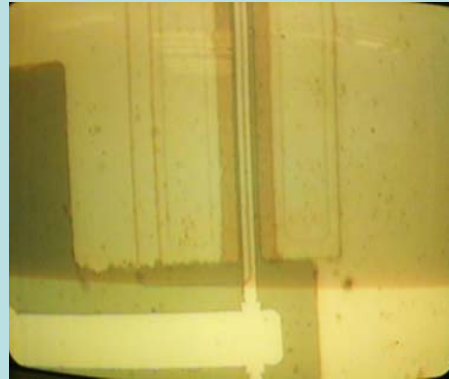


Elements for nanoelectronics

GaAs nano FETs



Structure



Optical microscopy image

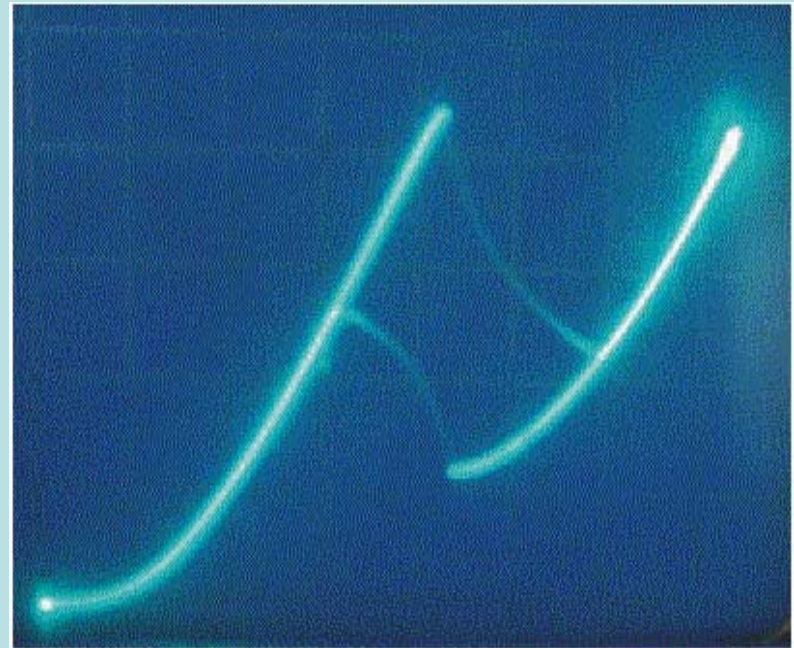
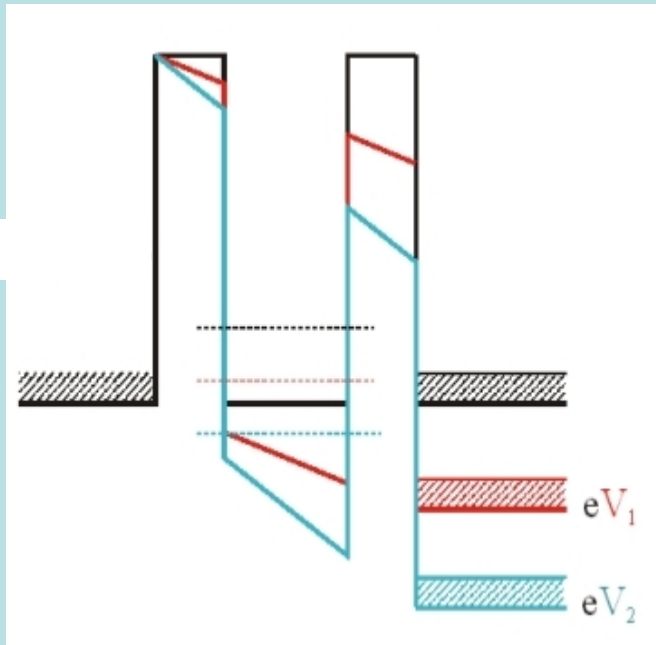


I-V characteristics

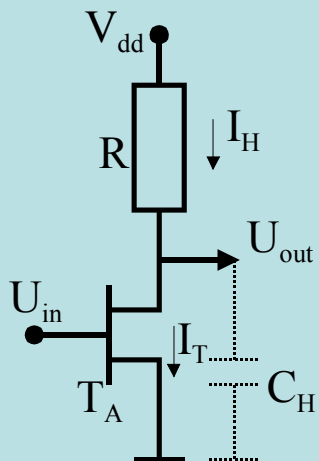


SPM gate image

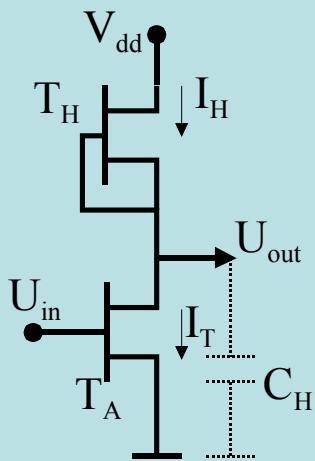
Resonant-tunneling diodes (RTD)



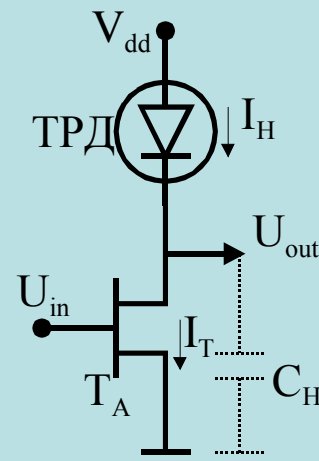
Electric circuits and current-voltage characteristics of basic MESFET inverters



Resistive load

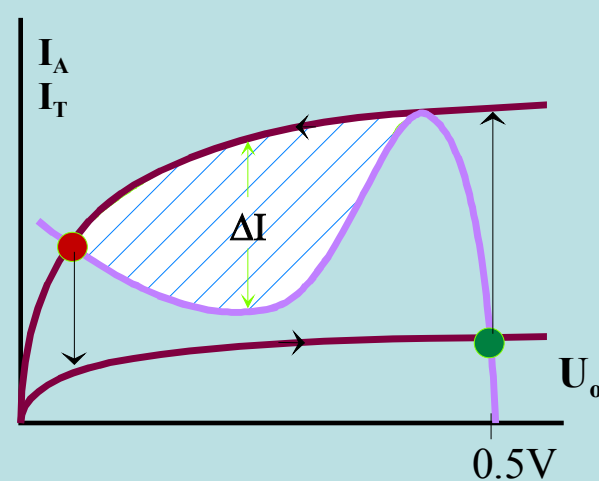
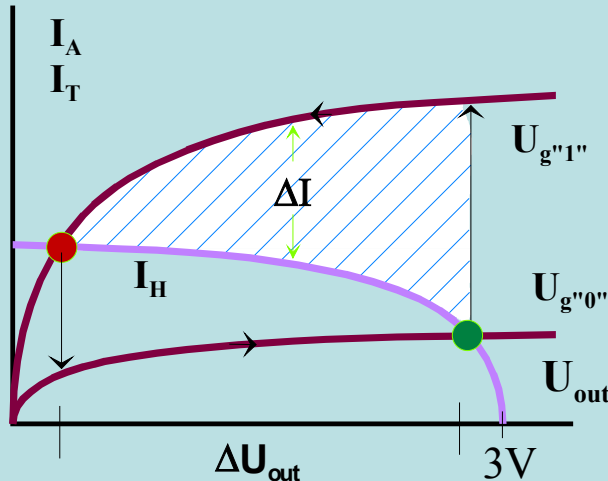
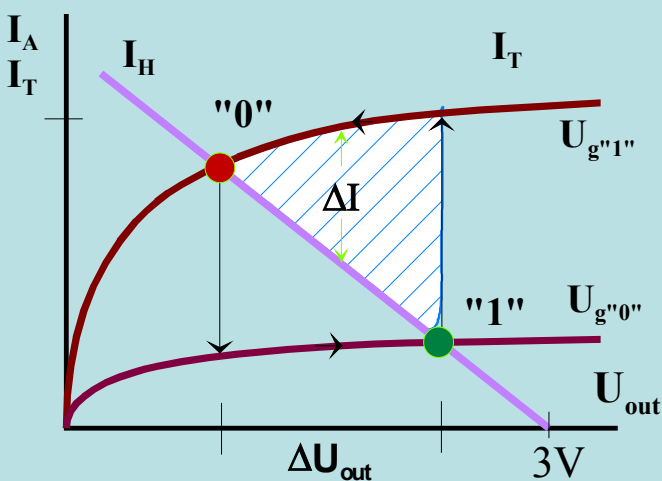


transistor load

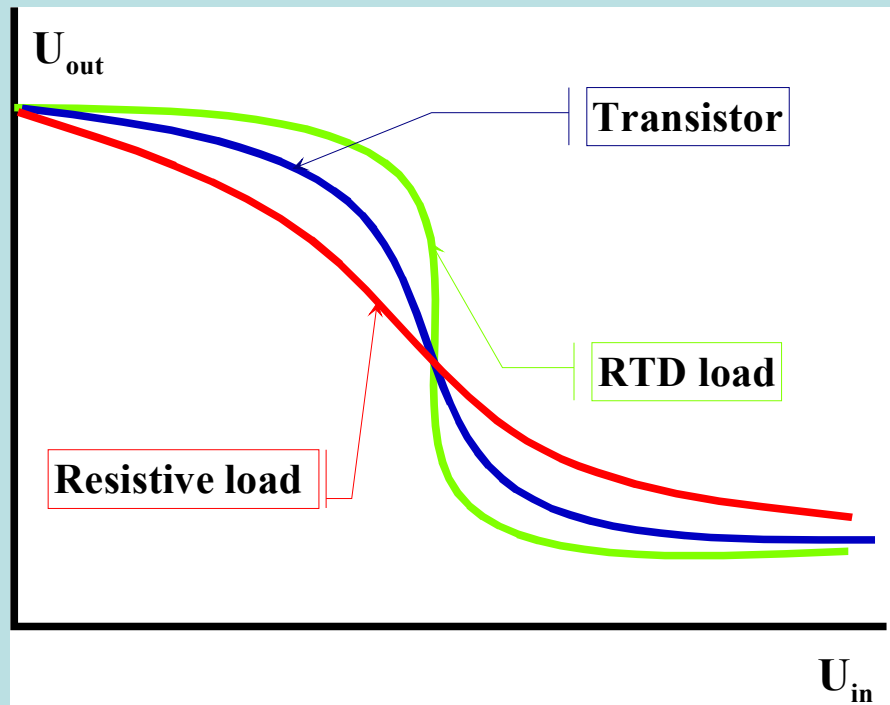


RTD as a load

Switching time $\Delta\tau = C_H \Delta U(I) / \Delta I$, for the RTD at the moment of switching C_H decreases due to the RTD negative capacitance.



Transfer characteristics of inverters



RTD:

1. High speed – up to tens of gigahertz with lithographic dimensions 0.6-0.8 μm
2. High gain
3. High interference margin
4. Power supply less than 1.5 V
5. Low power consumption

Number of elements necessary to construct logic functions based on different element base types

Circuit	TTL	CMOS	ECL	RTD/ MESFET
Two-stable-state XOR	33	16	11	4
Two-stable-state majority gate	36	18	29	5
Memory element (with 9 states)	24	24	24	5
2NO-OR+trigger	14	12	33	4
2NO – AND + trigger	14	12	33	4

Planarization problem



PERGAMON

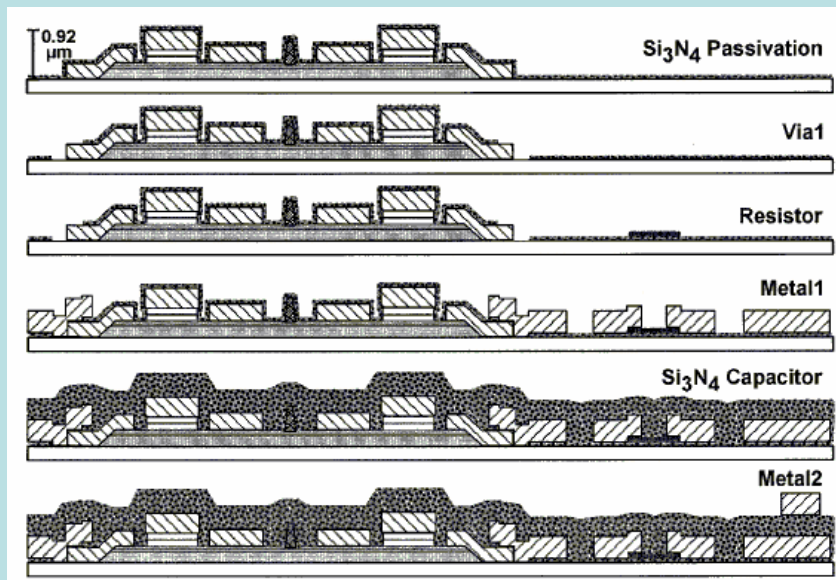
Solid-State Electronics 43 (1999) 1355-1365

**SOLID-STATE
ELECTRONICS**

Resonant-tunneling mixed-signal circuit technology

A. Seabaugh*, B. Brar, T. Broekaert, F. Morris, P. van der Wagt¹,
G. Frazier

Raytheon Systems Company, P.O. Box 660246, MS35, Dallas, TX 75266, USA



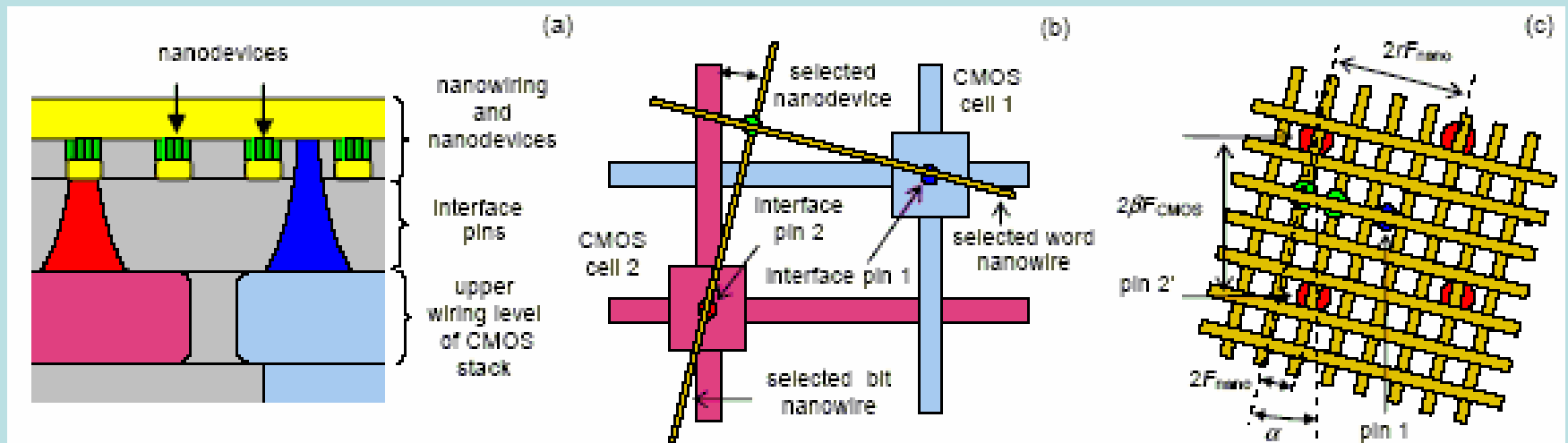
I-V curve of planar RTD+transistor chip



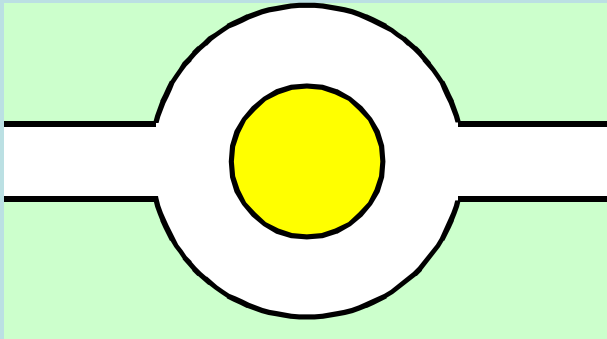
Transistor scaling limits – up to 22nm. And what is below?

**May be –
CMOL- electronics or
waveguide nanoelectronics!?**

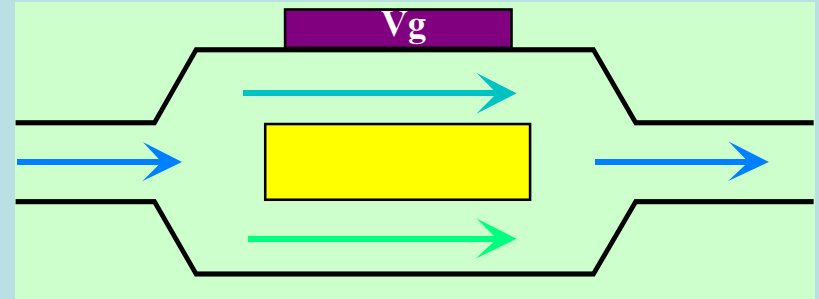
CMOL - electronics



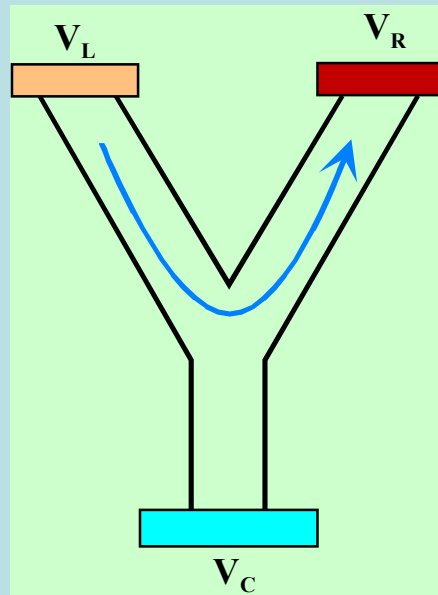
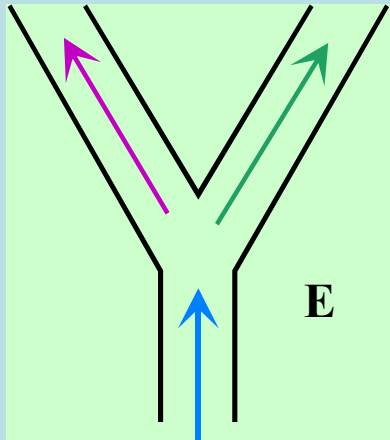
Quantum interference devices



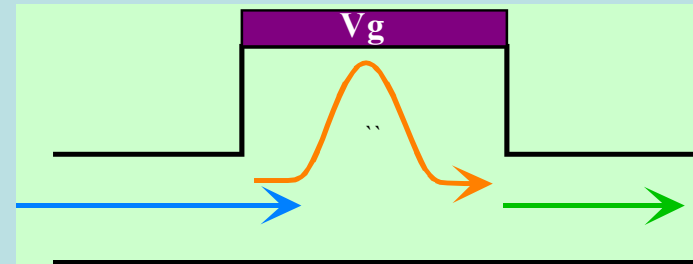
Aharonov-Bohm interferometer



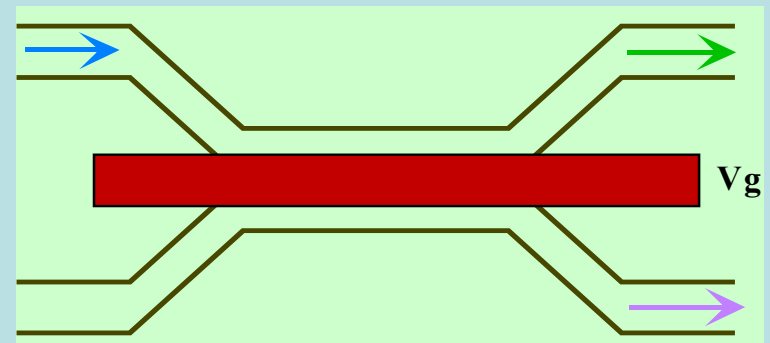
Quantum interference transistor



Three-terminal devices with Y-splitters



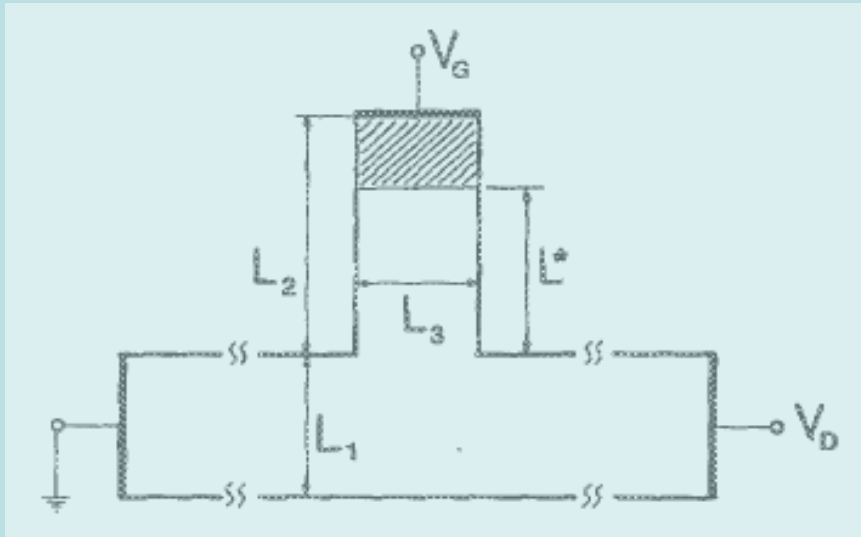
Quantum T-shaped transistor



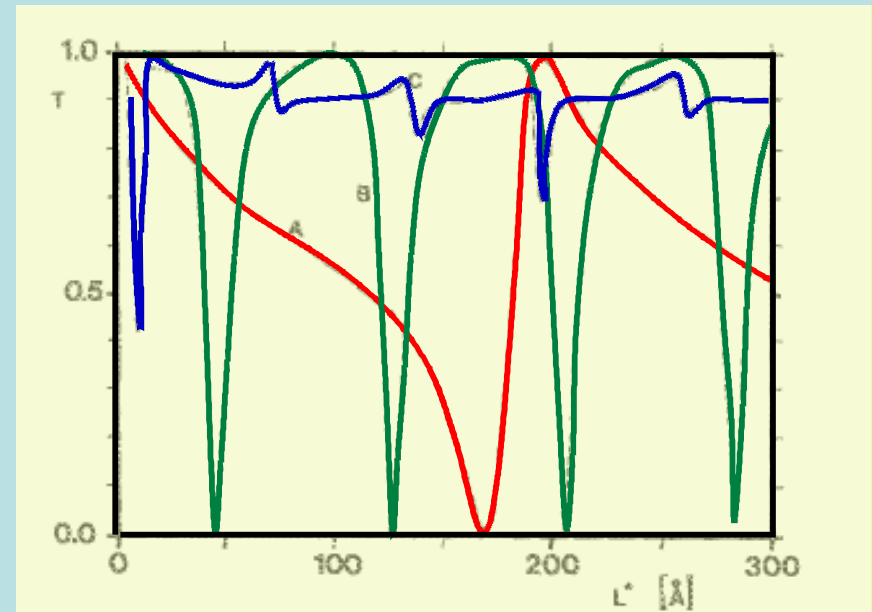
Directional splitter

T-shaped waveguide-like transistor

(F. Sols, M. Macucci, U. Ravaioli, K. Hess, Appl. Phys. Lett. 54 (4), 350, 1989)



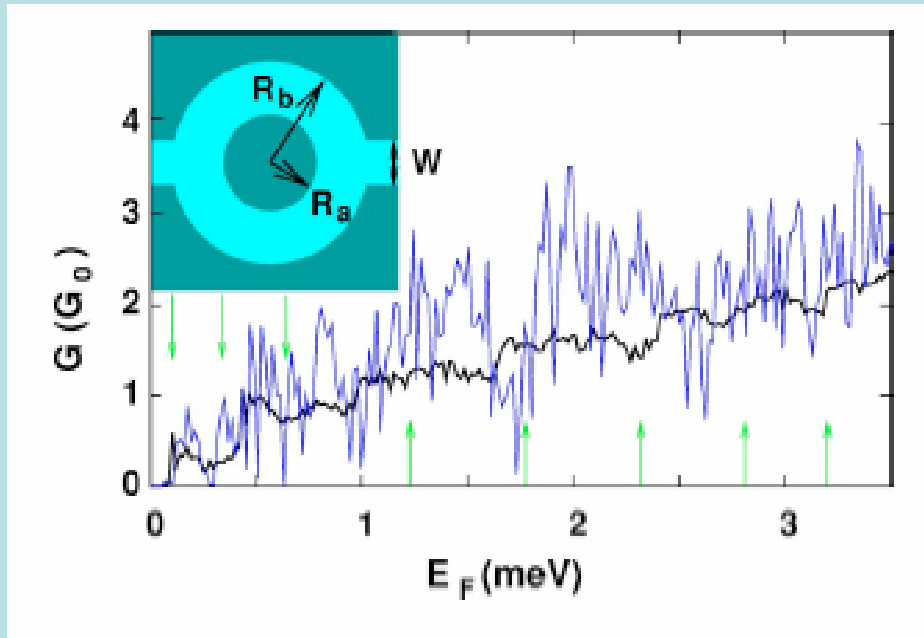
Schematic view of three-terminal structure



Transmission coefficient as function of the effective length of the stub for energy

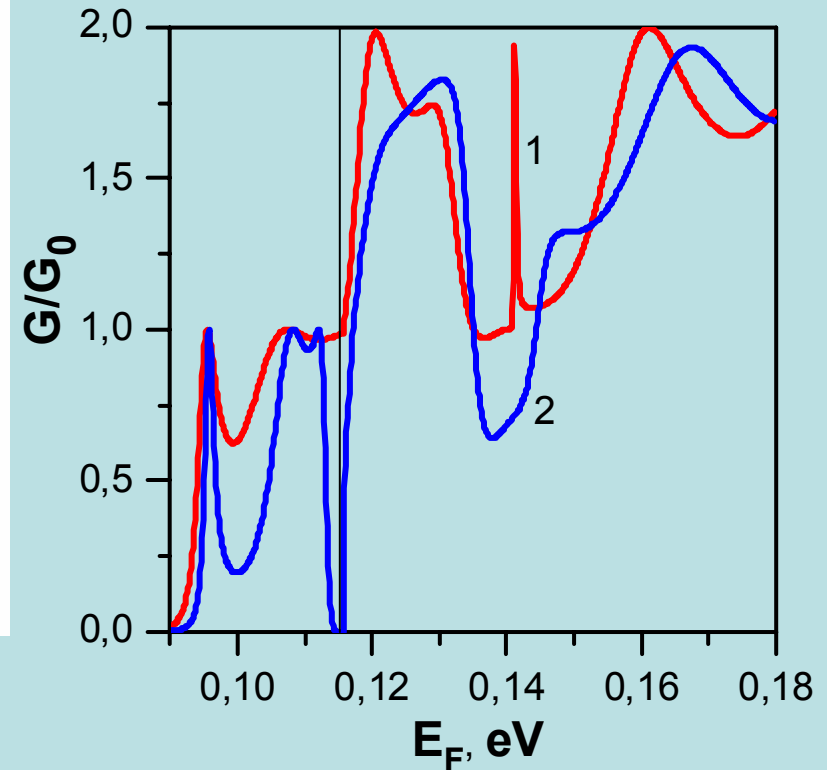
E (eV): A – 0.02; B – 0.118; C – 0.199

Complicated structure of conductance



Conductance of the AB ring as a function of the Fermi level of electrons. The arrows indicate the energies at which a new conducting channel in the lead is opened. $R_a = 350$ nm, $R_b = 630$ nm. $W = 200$ nm. $G_0 = 2e^2/h$

Phys. Rev. B **69**, 235304 (2004)

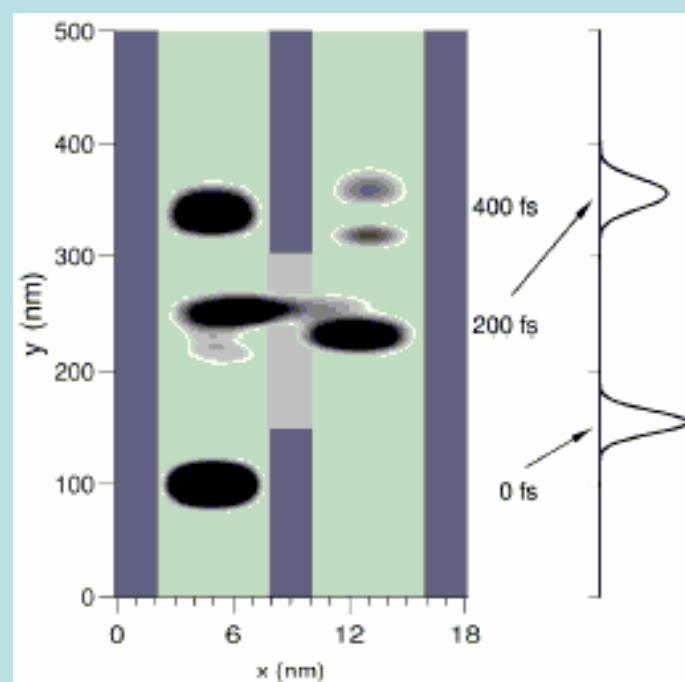
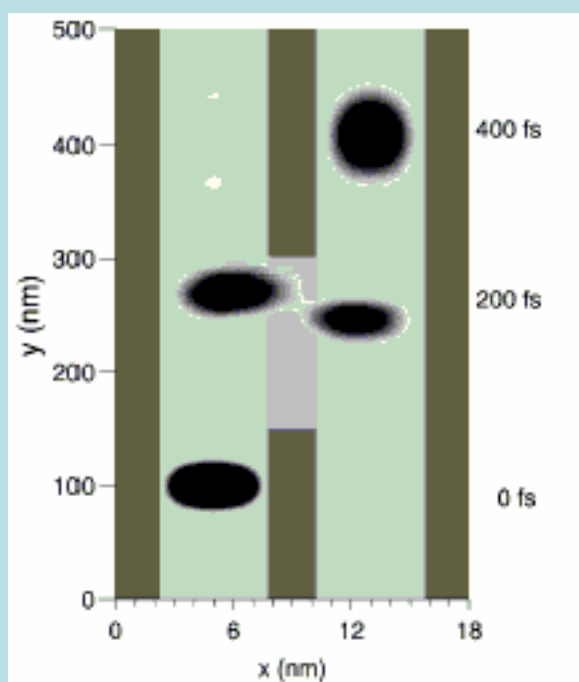


Quantum Logic Gates based on Coherent Electron Transport in Quantum Wires

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(Received 1 November 1999)

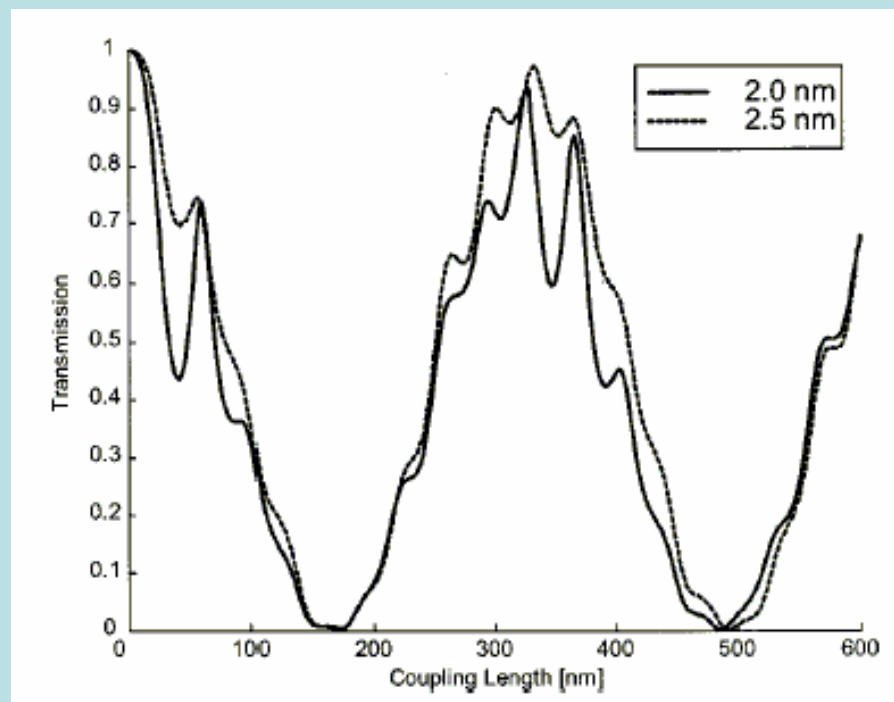
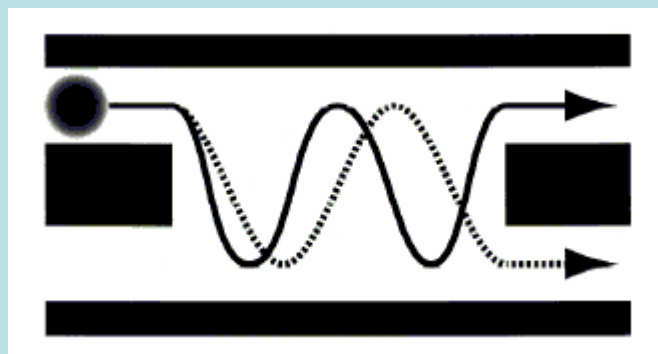


Magnetically switched quantum waveguide qubit

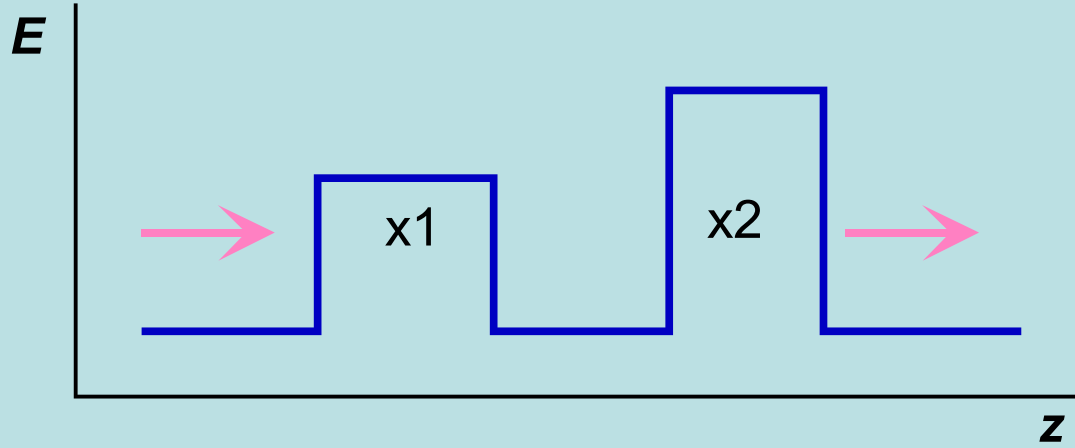
J. Harris, R. Akis, and D. K. Ferry^{a)}

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Arizona State University, Tempe, Arizona 85287-5706*

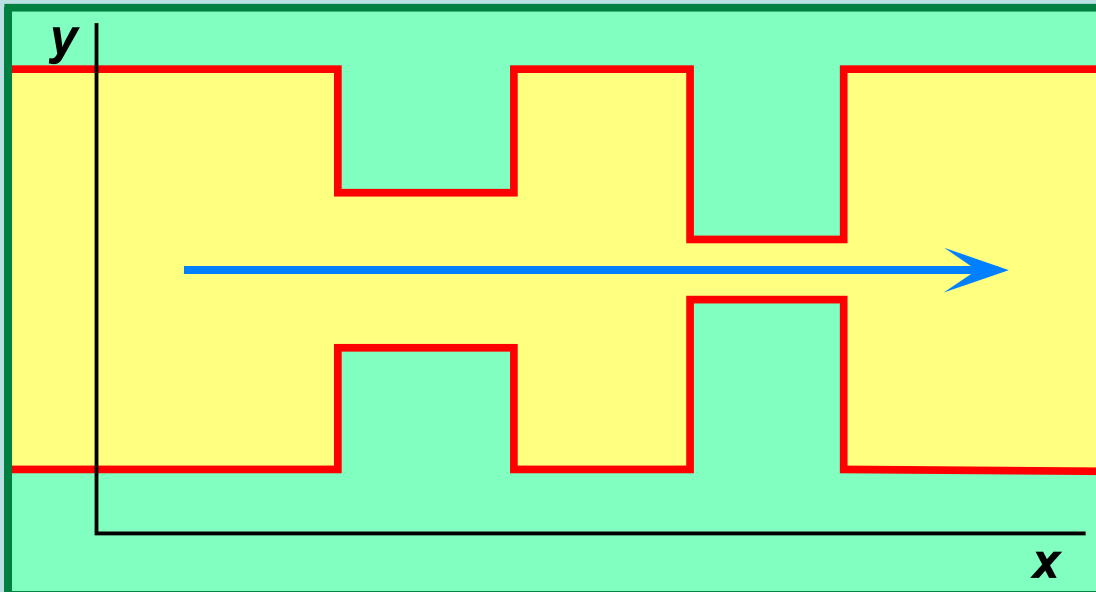
(Received 25 May 2001; accepted for publication 2 August 2001)



Analogy between electronic transport in semiconductor heterostructures and quantum wires with variable cross-section



Quantum Wells and Quantum Barriers on Heterostructure Potential
(GaAs/ $\text{Al}_x\text{Ga}_{1-x}\text{As}$) –
Energy Space

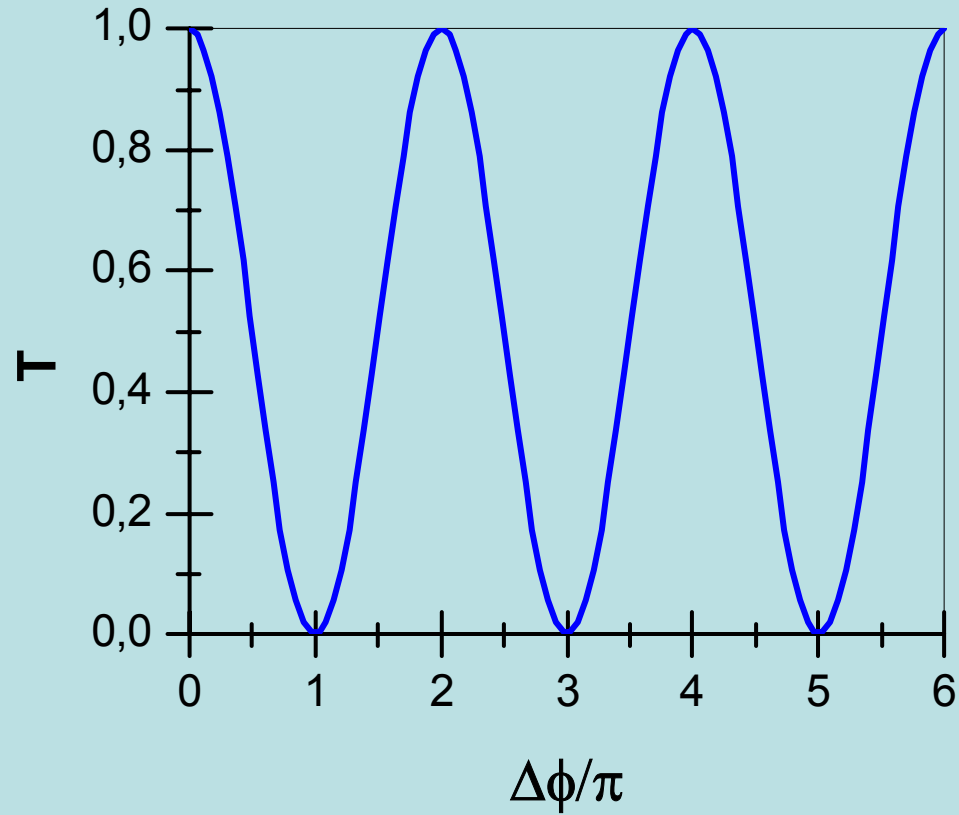


Electronic Waveguide –
Real Space
Gated 2D electron gas

Four basic resonances

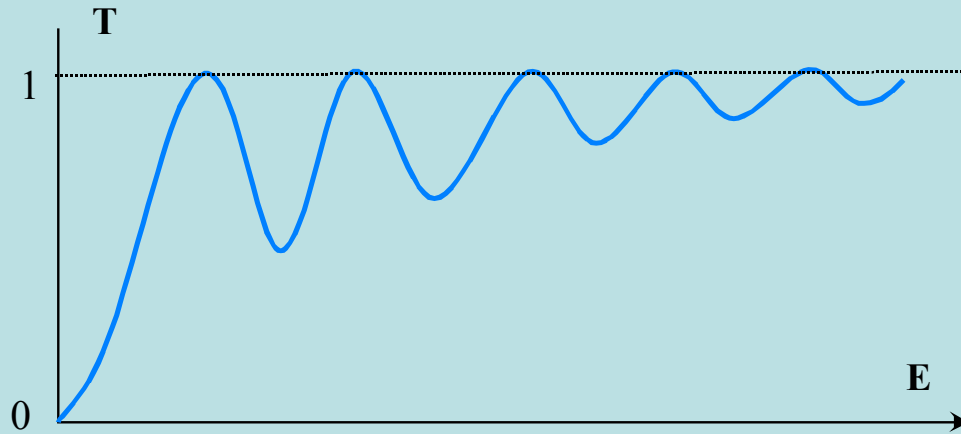
- **Resonant tunneling resonance**
- **Over-barrier (Ramsauer-Townsend-like) resonance**
- **Fano resonance**
- **Aharonov-Bohm resonance**

Aharonov-Bohm interference



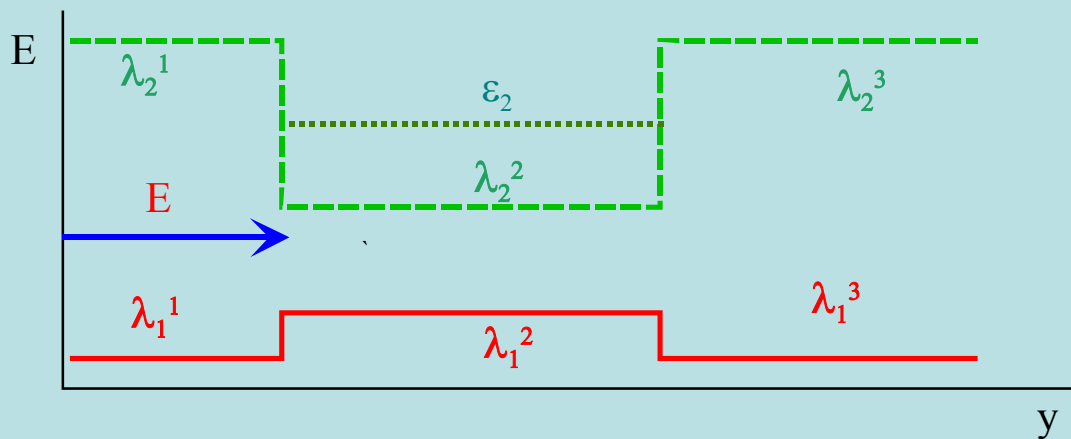
$$\Delta\phi = k_2 L_2 - k_1 L_1$$

Over-barrier resonances (Ramsauer-Townsend-like resonances)

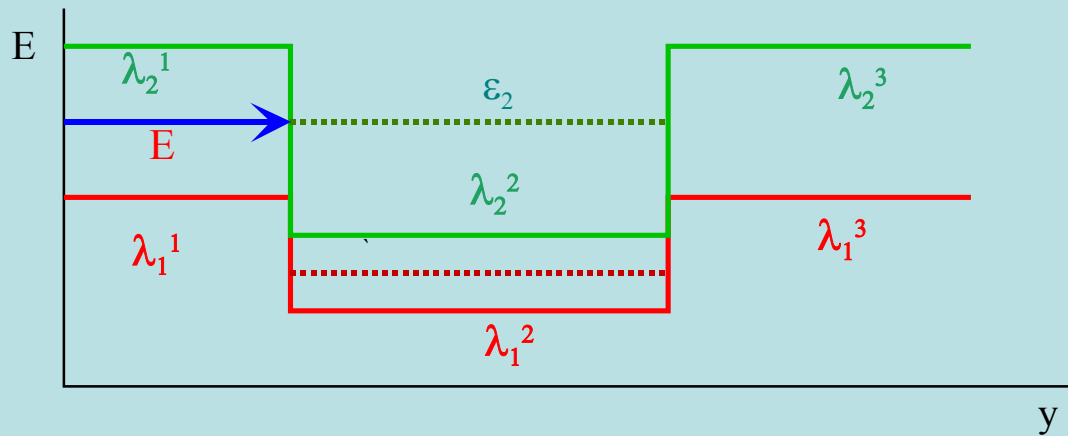
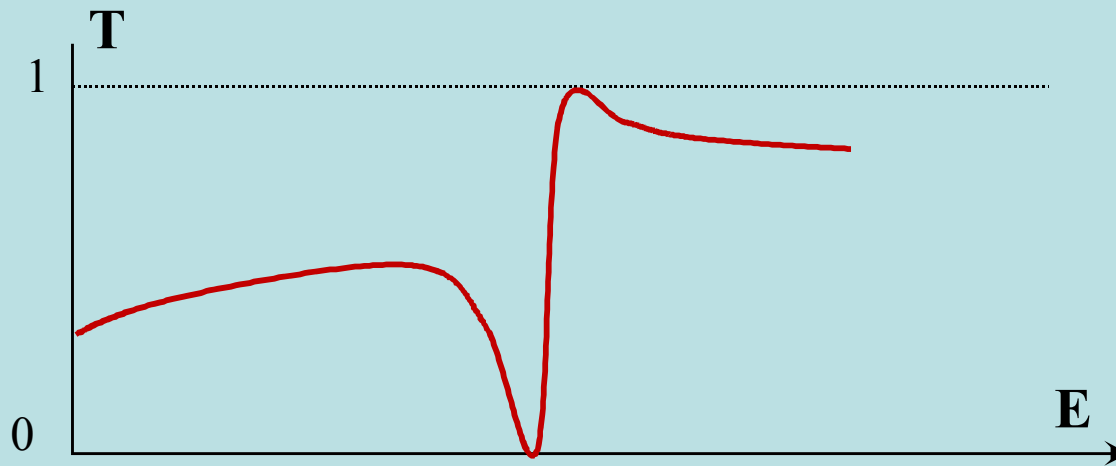


$$T = 1 \text{ for } \kappa_n^{(2)} d = n\pi$$

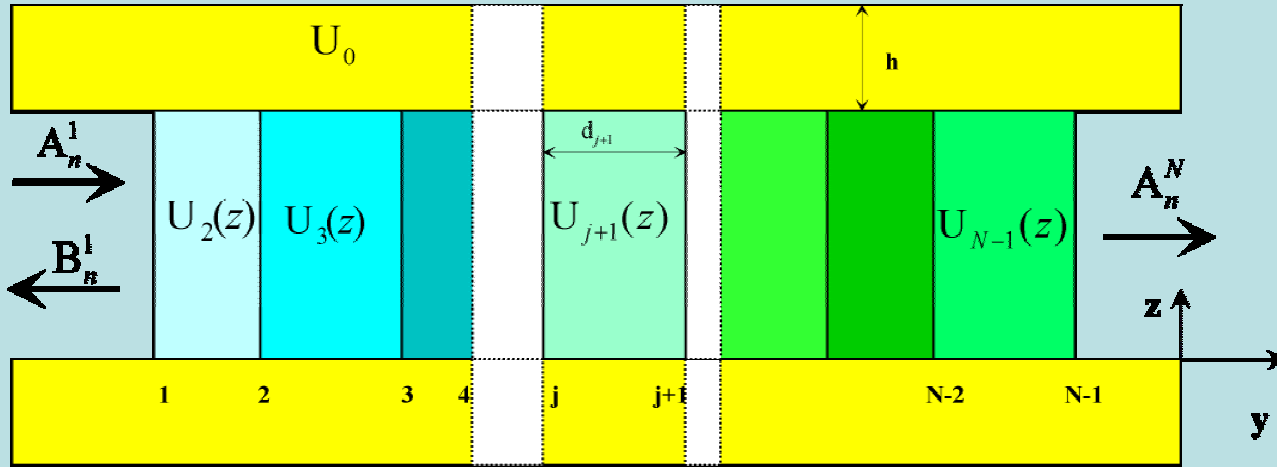
$$\text{where } \kappa_n^{(2)} = \sqrt{\frac{2m}{\hbar^2} (E - \lambda_n^2)}$$



Fano resonance



Solution of a scattering problem for a two-dimensional Schrodinger equation by expansion on waveguide modes



$$\frac{\partial^2 \psi(y, z)}{\partial y^2} + \frac{\partial^2 \psi(y, z)}{\partial z^2} + \frac{2m}{\hbar^2} (E - U(y, z)) \psi(y, z) = 0$$

$$\psi_j(y, z) = \sum_n \{A_n^j \exp[i\kappa_n^j (y - y_j)] + \{B_n^j \exp[-i\kappa_n^j (y - y_j)]\} Z_n^j(z)$$

$$\kappa_n^j = \sqrt{\frac{2m}{\hbar^2} (E - \lambda_n^j)} , \quad \lambda_n^j$$

- eigenvalues of the one-dimensional Schrodinger equations with a potential $U_j(z)$, $Z_n^j(z)$ a - corresponding wave functions.

From a continuity condition at $y = y_j$ for ψ and $\frac{\partial \psi}{\partial y}$ we obtain relations

$$A_n^{j+1} = \frac{1}{2} \sum_m \left\{ \left(1 + \frac{\kappa_m^j}{\kappa_n^{j+1}}\right) \mu_{mn} \exp\{i\kappa_m^j d_j\} A_m^j + \left(1 - \frac{\kappa_m^j}{\kappa_n^{j+1}}\right) \mu_{mn} \exp\{-i\kappa_m^j d_j\} B_m^j \right\}$$

$$B_n^{j+1} = \frac{1}{2} \sum_m \left\{ \left(1 - \frac{\kappa_m^j}{\kappa_n^{j+1}}\right) \mu_{mn} \exp\{i\kappa_m^j d_j\} A_m^j + \left(1 + \frac{\kappa_m^j}{\kappa_n^{j+1}}\right) \mu_{mn} \exp\{-i\kappa_m^j d_j\} B_m^j \right\}$$

$$\mu_{mn} = \int Z_m^j(z) Z_n^{j+1}(z) dz$$

In the matrix form

$$\begin{pmatrix} A \\ B \end{pmatrix}^{j+1} = D_j \begin{pmatrix} A \\ B \end{pmatrix}^j$$

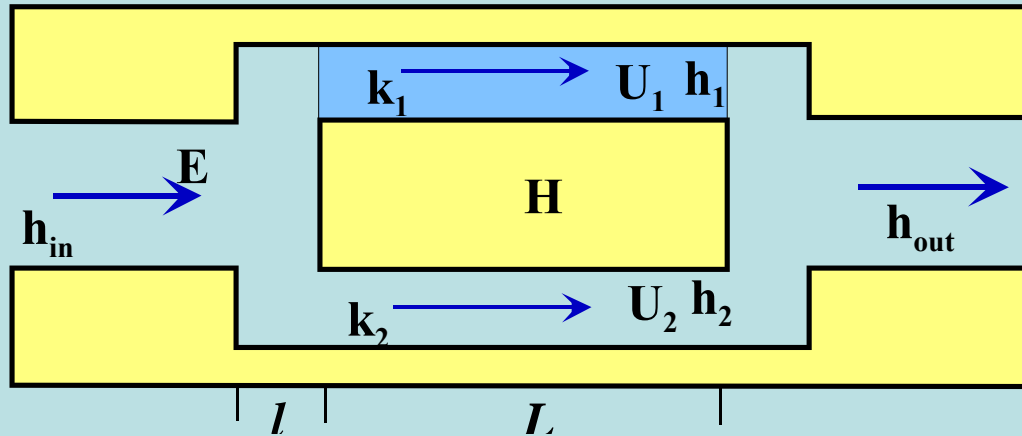
Denoting $A^1 \equiv A$; $B^1 \equiv r$; $A^N \equiv t$ and sequentially applying, we obtain

$$\begin{pmatrix} t \\ 0 \end{pmatrix} = D \begin{pmatrix} A \\ r \end{pmatrix}$$

The **transfer matrix** of the structure

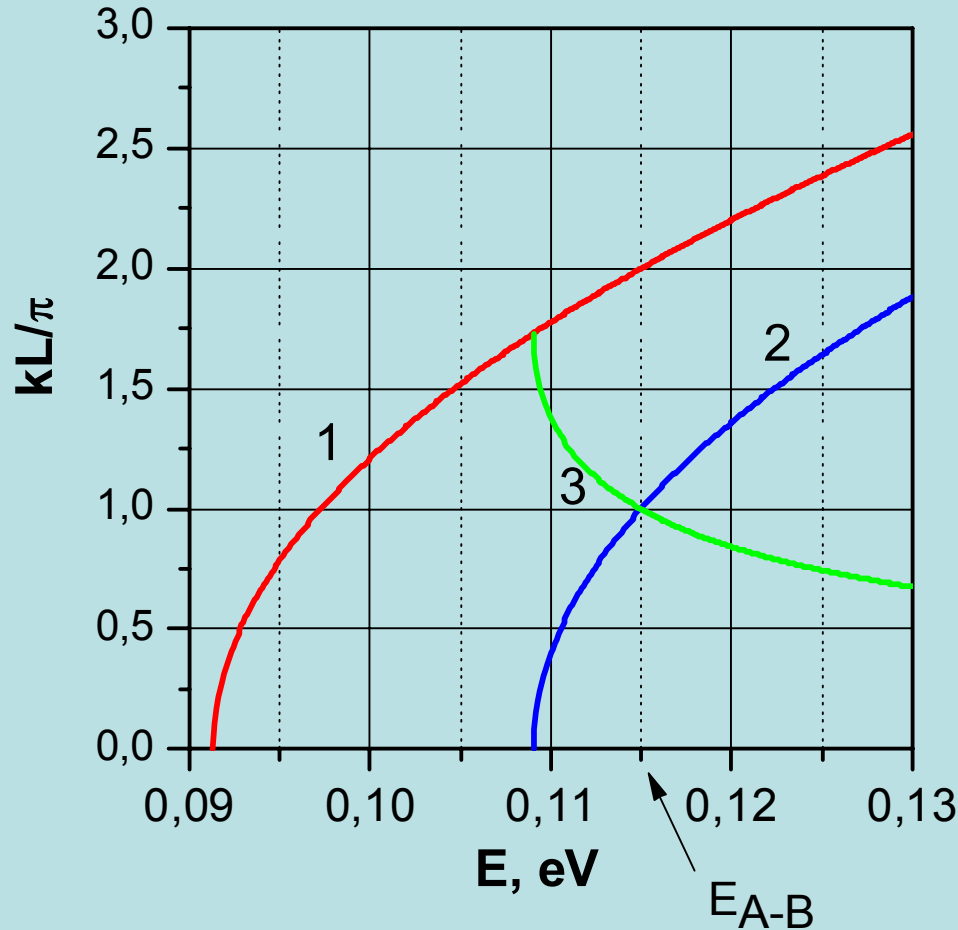
$$D = D_{N-1} D_{N-2} \cdots D_2 D_1$$

Geometry of Aharonov-Bohm interferometer



Step-like transition region

Optimal Channel Length



at $k_2 L_2 = \pi$ $k_1 L_1 = 2\pi$

**Over-barrier
resonance in each
channel and total
minimum due to A-B
interference**

$h_1 = h_2 = 50$ A, $L = 309$ A

$E_{A-B} = 0.115$ eV

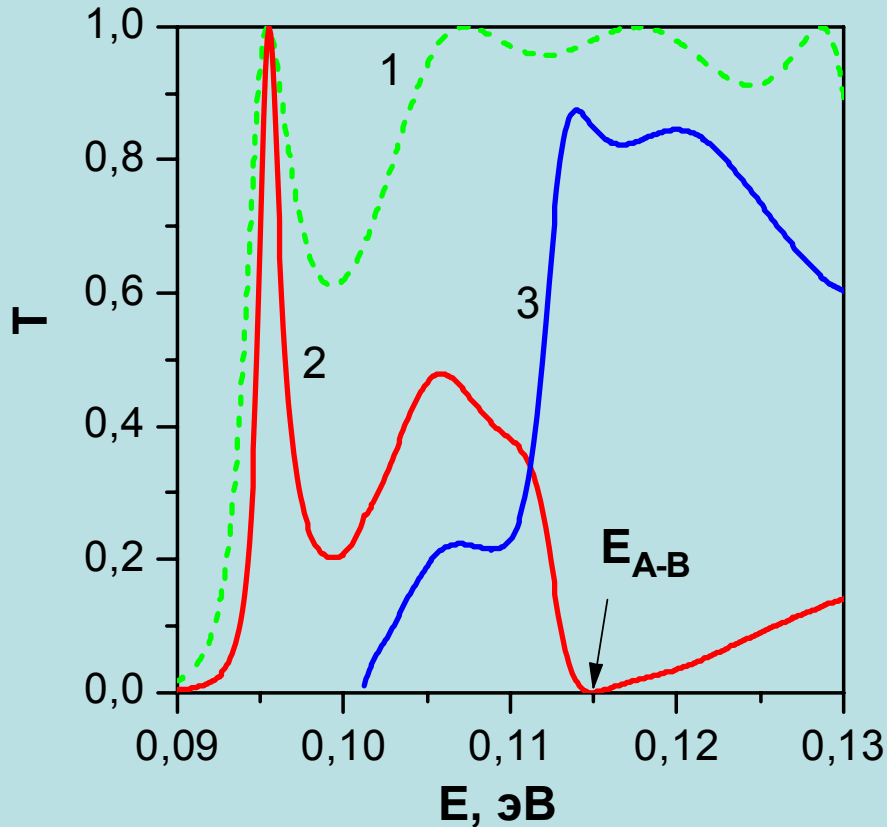
$\Delta U = 0.02$ eV

Phase change in channels 1 and 2 and interchannel phase difference (3)

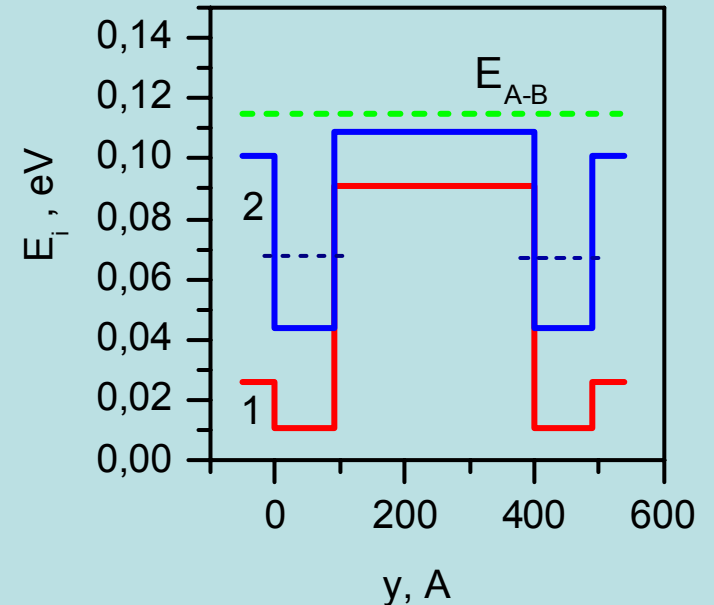
Choice of incoming wave-guide widths

$$h_{in} = h_{out} = 120 \text{ \AA}$$

Absence of Fano resonance $E_L^2 > E_{in}^2$



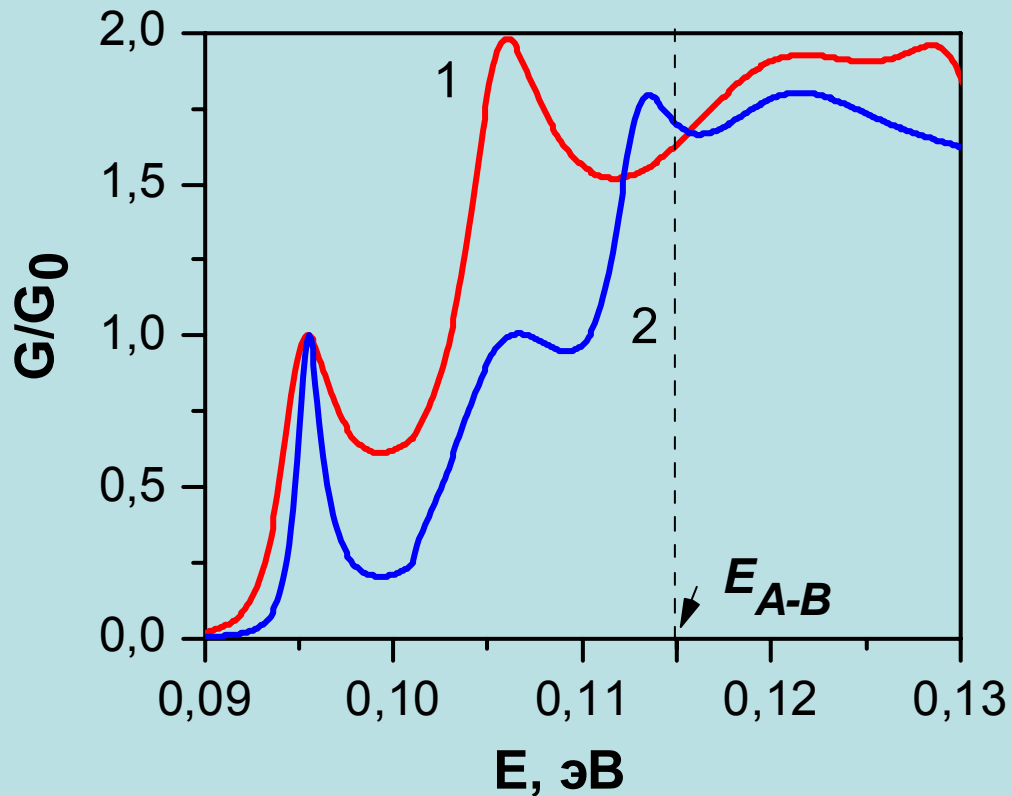
But $E_L^2 < E_{A-B}$



1-st mode transmission at $\Delta U=0$ (1)
 1-st (2) and 2-nd (3) modes transmission at $\Delta U=0.02 \text{ eV}$

Transverse(z) quantization energy position along the direction of the wave-guide (y)

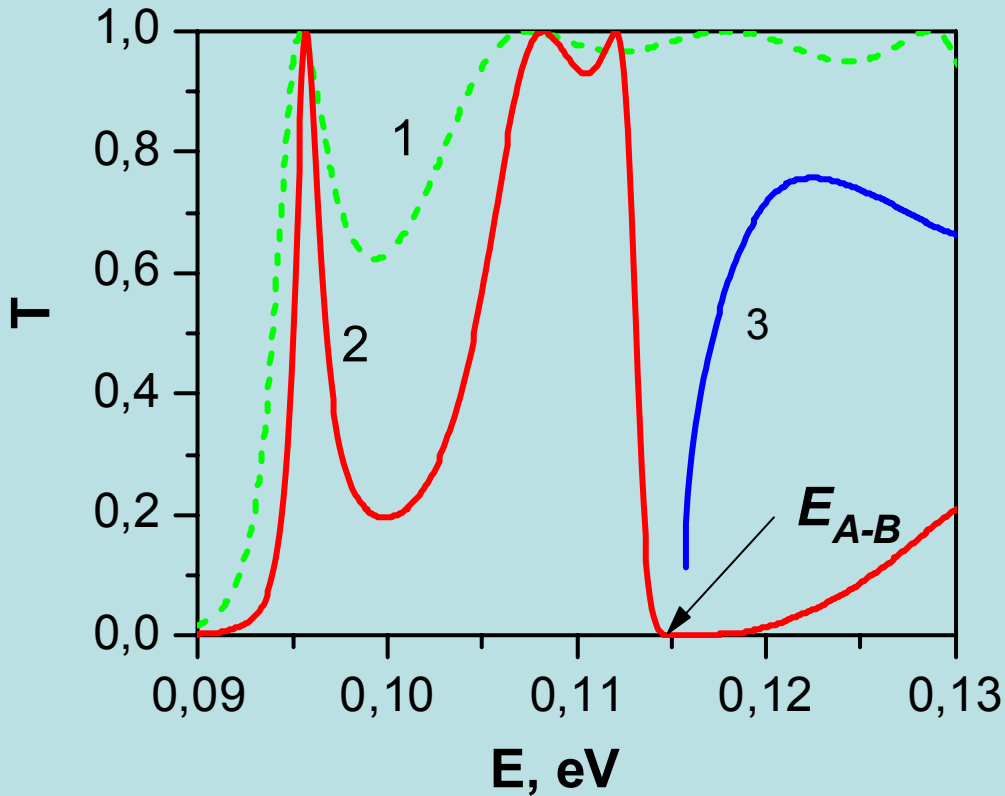
Conductance



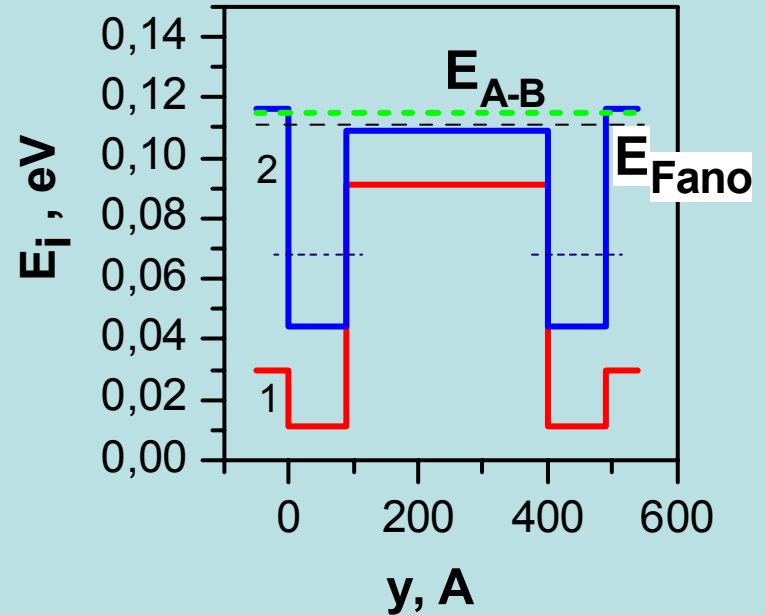
Conductance dependence on Fermi energy for $h_{in} = 120$ A

1 - $\Delta U = 0$, 2 - $\Delta U = 0.02$ eV

Transmittance for $h_{in}=110$ A

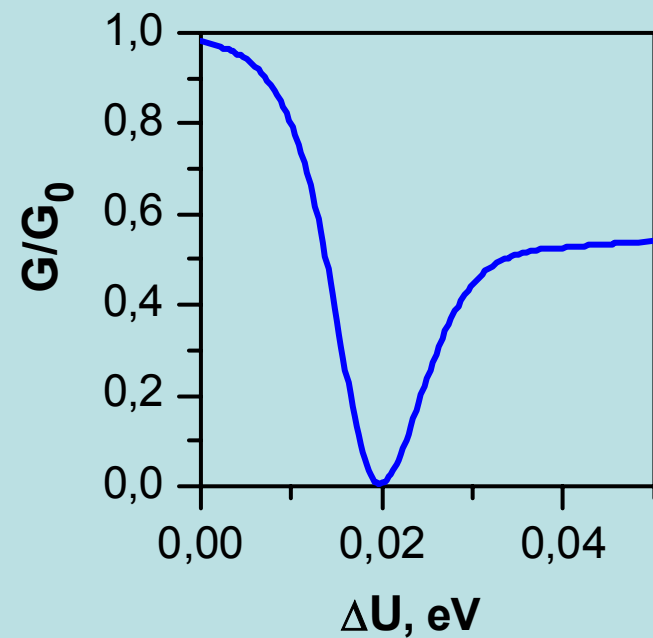
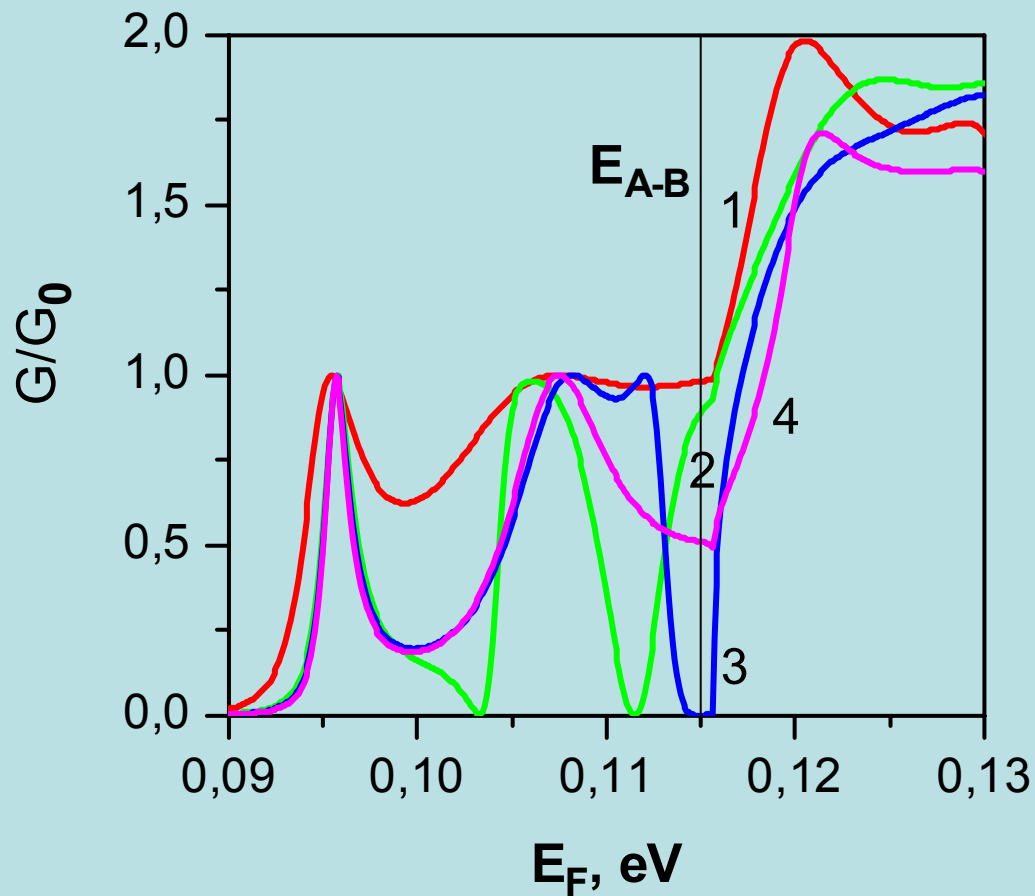


Transmission in 1st mode for $\Delta U=0$ (1) and in 1st (2) and 2nd (3) modes for $\Delta U=0.02$ eV



Quantization energy profile along the direction of the waveguide (y)

Variation of conductance with changing of ΔU for $h_{in}=110 \text{ A}$



$G(\Delta U)$ curve at $E_F = 0.115 \text{ eV}$

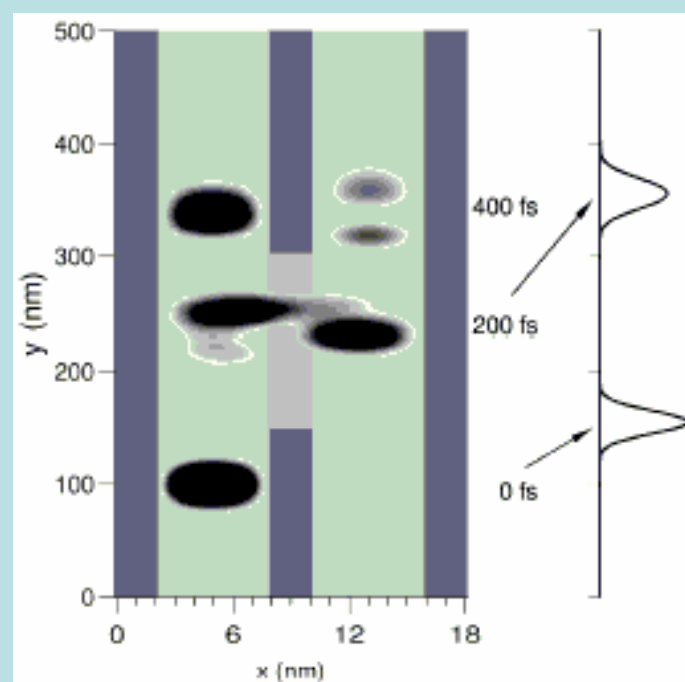
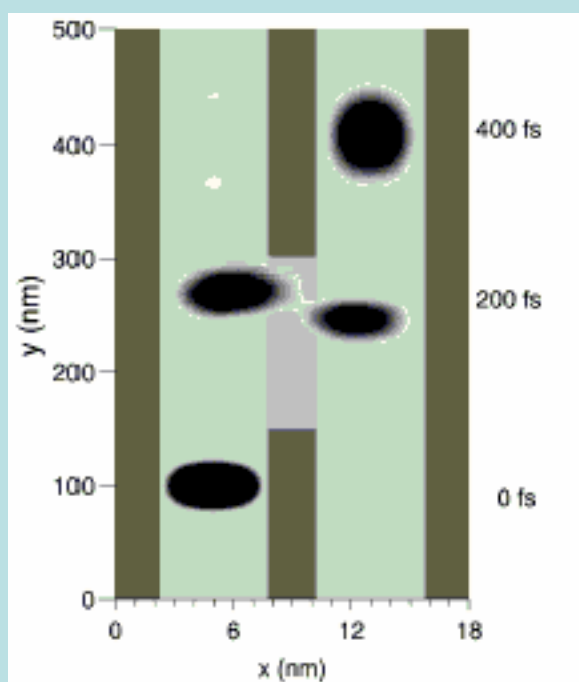
1 ΔU (meV): 1 – 0; 2 – 10; 3 – 20; 4 – 50

Quantum Logic Gates based on Coherent Electron Transport in Quantum Wires

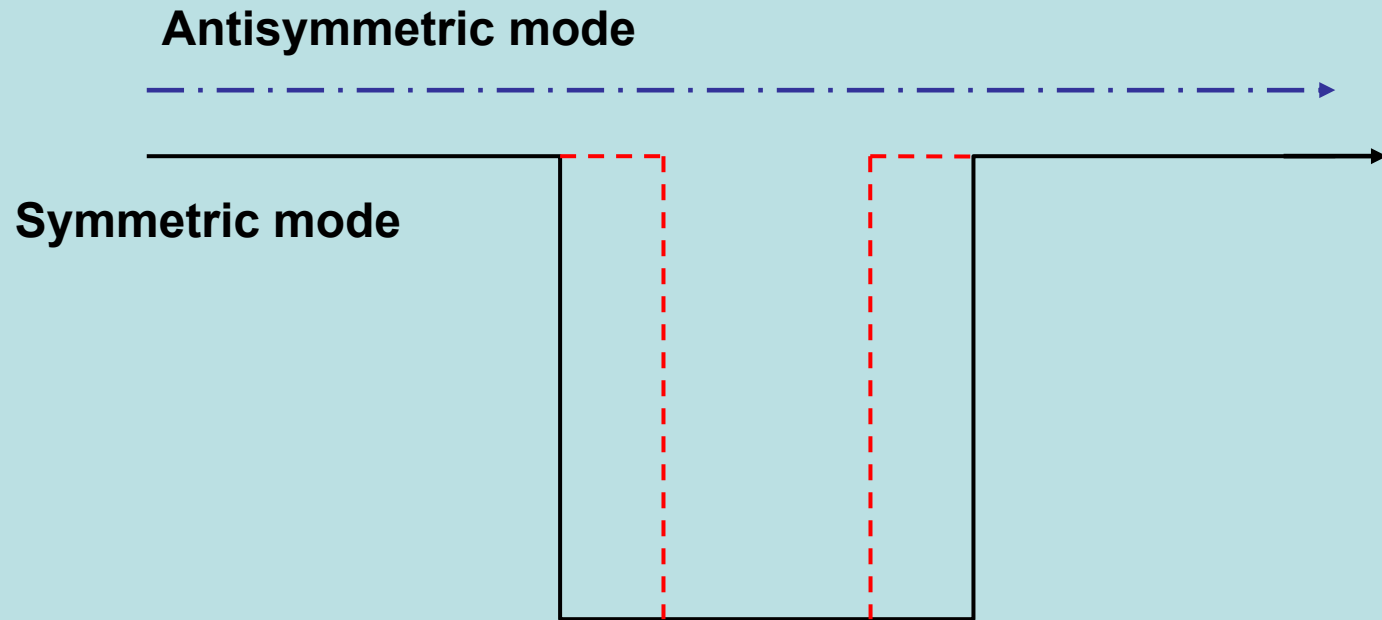
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The role of a window in a waveguide barrier

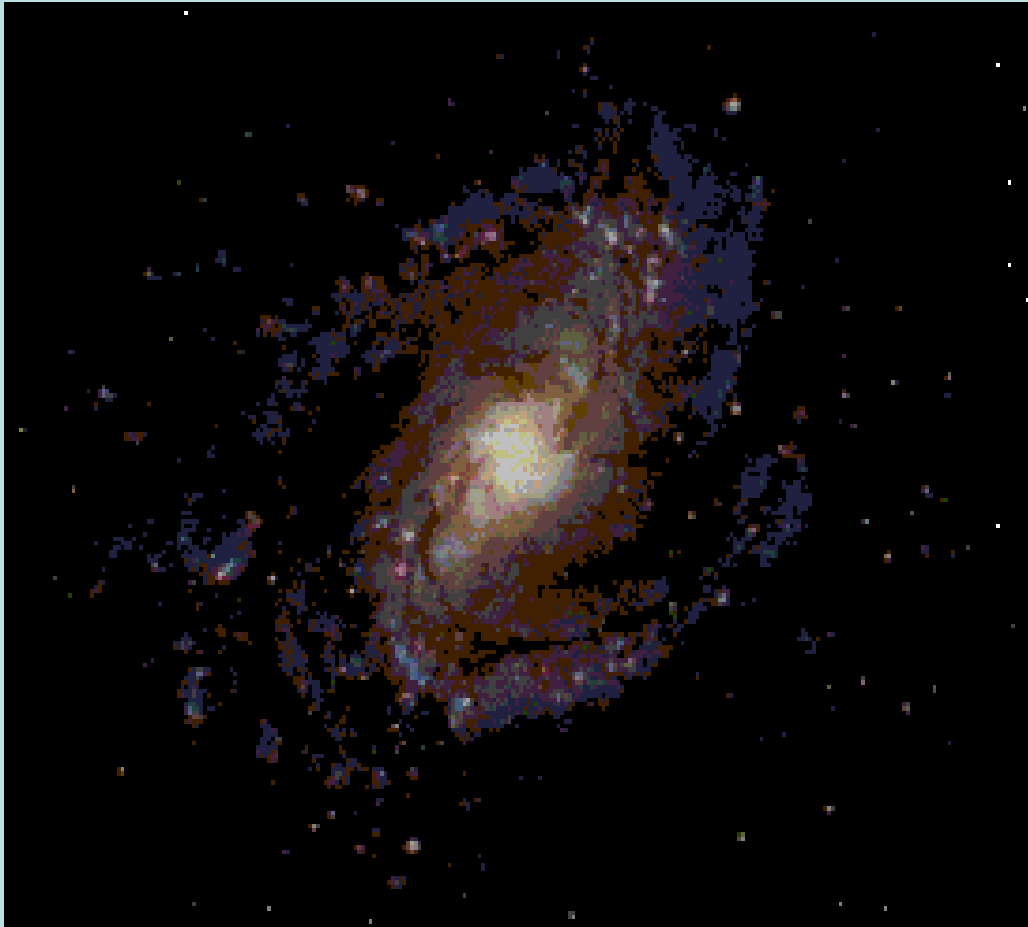


Fundamental problems of physics in low dimensions

The fractional quantum Hall effect ...is important ...primarily for one:...particles carrying an exact fraction of the electron charge e and ...gauge forces between these particles, two central postulates of the standard model of elementary particles, can arise spontaneously as **emergent** phenomena....idea whether the properties of the universe ...are **fundamental** or **emergent**...I believe...must give string theorists pause...

R.B.Laughlin “Nobel Lecture: Fractional quantization” Rev. Mod. Phys. v. 71, p. 863 (1999)

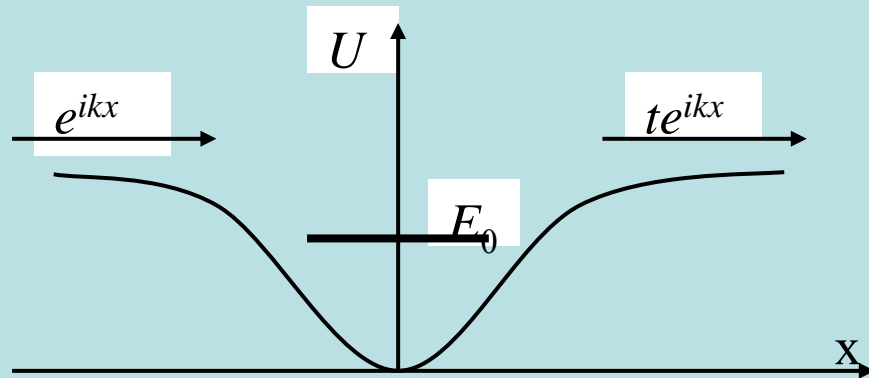
Dark matter



At present it is established that from 88% to 95% of matter in the Universe are of unknown origin (neutrino mass, new particles etc.?)

Can objects exist in crystal
which are invisible?

Reflectionless quantum mechanical potentials



$$U(x) = -U_0 c \hbar^{-2} \alpha x$$

$$1 + (8mU / \hbar^2 \alpha^2) = (2n + 1)^2$$

$$r \equiv 0, \quad t = e^{i\varphi(k)}$$

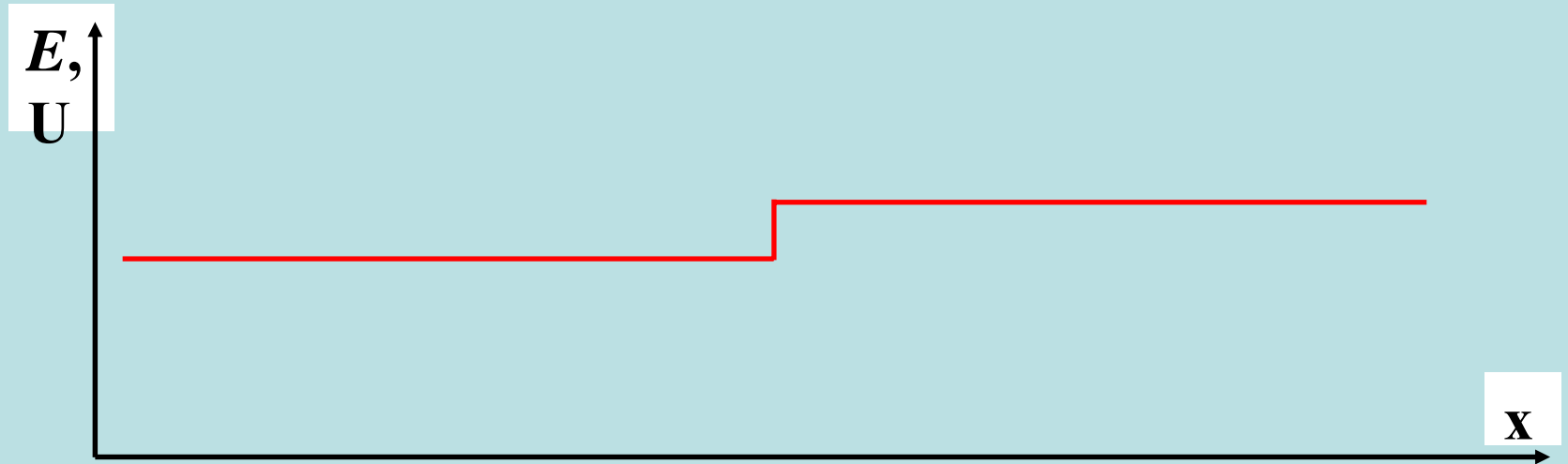
$$E(k) = \frac{\hbar^2 k^2}{2m}, \quad E_0 = -\frac{\hbar^2 \alpha^2}{2m}$$

$$\varphi = 2 \operatorname{arctg} \frac{\alpha}{k}$$

There **exist** *hidden microscopic objects* in crystals – extended objects which are “invisible” in respect to low-energy electron energy ($r \equiv 0$ и $t \equiv 1$ in continuum limit)

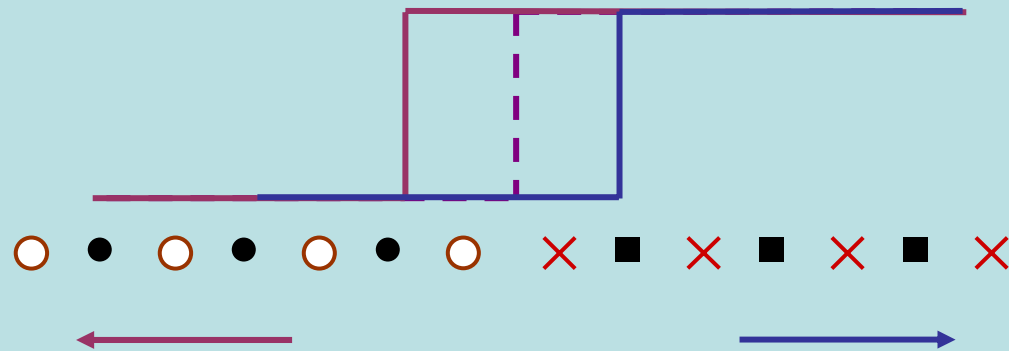
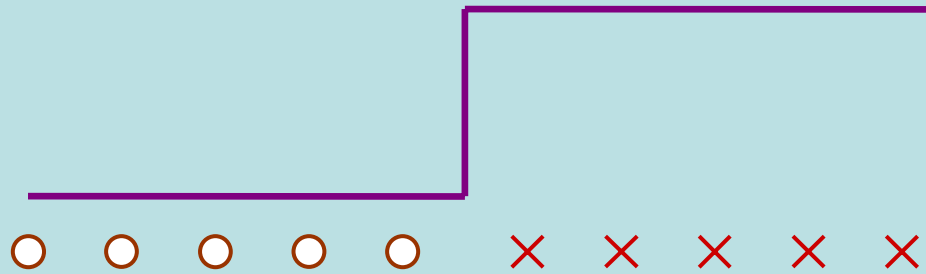
A.A.Gorbatsevich “Hidden Defect Pairs: Objects invisible in low energy electron scattering” e-print cond-mat/0511054 (2005)

Envelope Function (or Effective Mass) Method

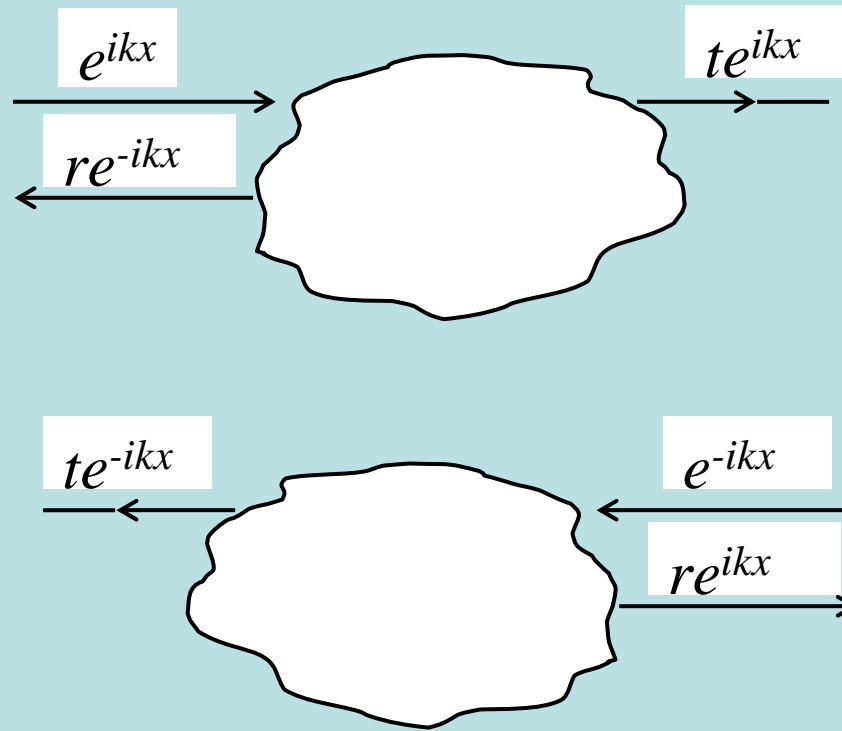


$$\left[-\frac{\hbar^2 \nabla^2}{2m^*} + U(\mathbf{r}) \right] \psi(\vec{r}) = E\psi(\vec{r})$$

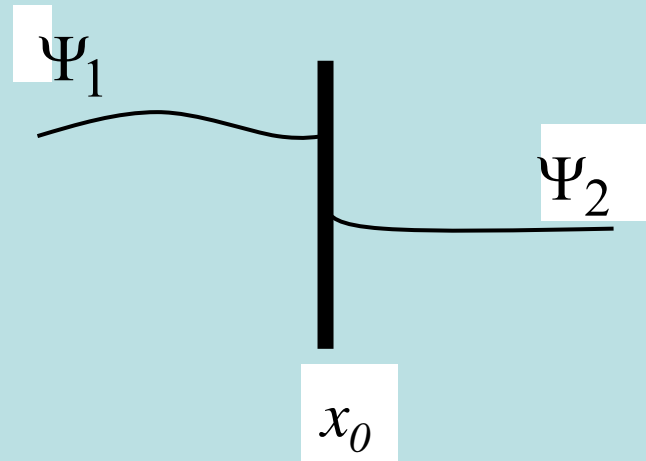
Location of heterointerface can be determined exactly!



Scattering data



Generalized Boundary Conditions



$$\begin{pmatrix} \Psi \\ \nabla\Psi \end{pmatrix}_{x_{0+}} = \hat{T} \begin{pmatrix} \Psi \\ \nabla\Psi \end{pmatrix}_{x_{0-}} = \begin{pmatrix} T_{11} & T_{12} \\ T_{21} & T_{22} \end{pmatrix} \begin{pmatrix} \Psi \\ \nabla\Psi \end{pmatrix}_{x_{0-}}$$

Extraction procedure

$$\psi = \begin{cases} e^{ik_1(x-\delta_1)} + re^{-ik_1(x-\delta_1)}, & x < x_0. \\ te^{-ik_2(x-\delta_2)}, & x > x_0. \end{cases}$$

$$r = e^{2ik_1(x_0-\delta_1)} \frac{k_1 T_{22} - k_2 T_{11} - iT_{21} - ik_1 k_2 T_{12}}{k_1 T_{22} + k_2 T_{11} + iT_{21} - ik_1 k_2 T_{12}}$$

$$t = e^{i[k_1(x_0-\delta_1)-k_2(x_0-\delta_2)]} \sqrt{\frac{m_2}{m_1}} \frac{2k_1}{k_1 T_{22} + k_2 T_{11} + iT_{21} - ik_1 k_2 T_{12}}$$

$$r \approx e^{2ik_1Na} \left[1 - 2 \frac{k_2}{k_1} T_{11}^2 - 4ik_2 T_{11}^2 (\Delta x_0 - \delta_1) + \dots \right],$$

$$r' \approx -e^{-2ik_2Na} \left[1 - 2 \frac{k_2}{k_1} T_{11}^2 - 2ik_2 (\Delta x_0 - \delta_2 - T_{12} T_{11}) + \dots \right],$$

$$t^{-1} \approx \frac{1}{2} e^{i(k_2 - k_1)Na} \sqrt{\frac{m_1}{m_2}} T_{22} \left[1 + \frac{k_2}{k_1} T_{11}^2 - ik_2 [(T_{11} T_{12} + T_{11}^2 (\Delta x_0 - \delta_1) - \Delta x_0 + \delta_2)] + \dots \right]$$

Model



$$\vec{r} \rightarrow -\vec{r}$$

$$\hat{H} = \sum_{j=1,2} \sum_{n_1, n_2} \left[\varepsilon_j C_{j n_j}^+ C_{j n_j} - C_{1 n_1}^+ (t_+ C_{2 n_1+1} + \right.$$

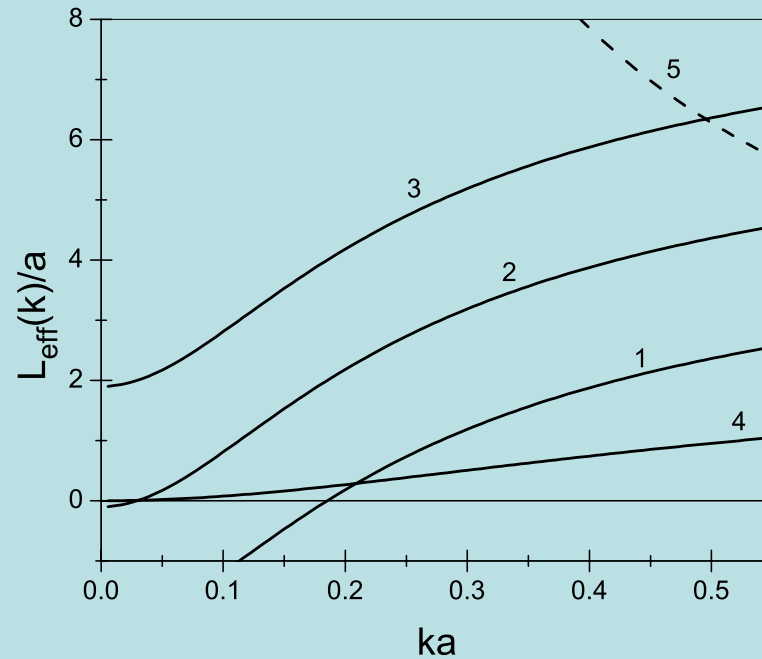
$$\left. + t_- C_{2 n_1-1}) + \varepsilon_j^* C_{M_j}^+ C_{M_j} + h.c. \right]$$

Scattering data in the model

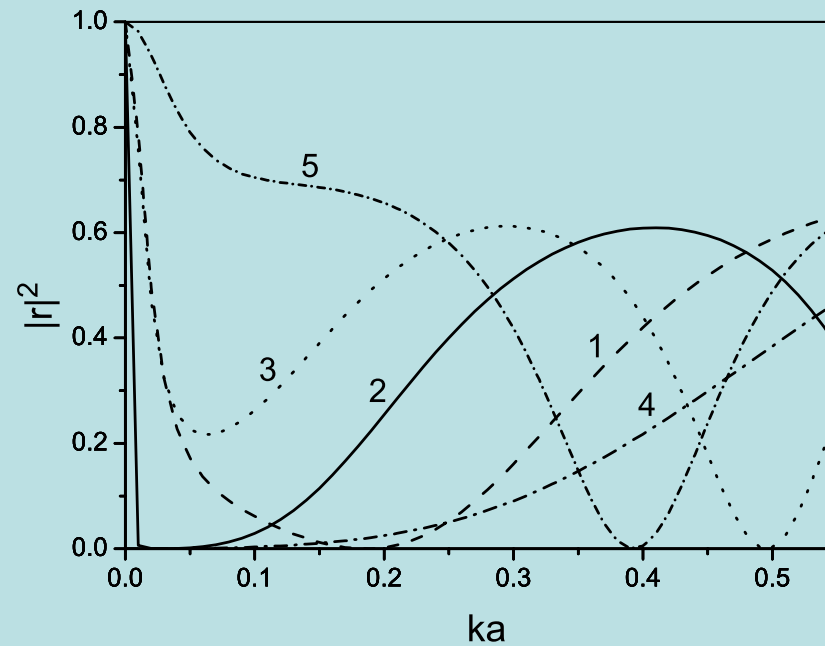
$$r = -e^{2ikM_1 - 2i\zeta_a} \frac{(\Delta_1 \Delta_2 u_k v_k + i\Delta_{k-} U_k) \sin[L_{eff}(k)k]}{\Delta_1 \Delta_2 u_k v_k \sin[L_{eff}(k)k] - (\Delta_{k+} + 2iu_k v_k U_k) U_k e^{-iL_{eff}(k)k}}$$

$$t = -e^{-iL_{eff}(k)k} \frac{2iu_k v_k U_k^2}{\Delta_1 \Delta_2 u_k v_k \sin[L_{eff}(k)k] - (\Delta_{k+} + 2iu_k v_k U_k) U_k e^{-iL_{eff}(k)k}}$$

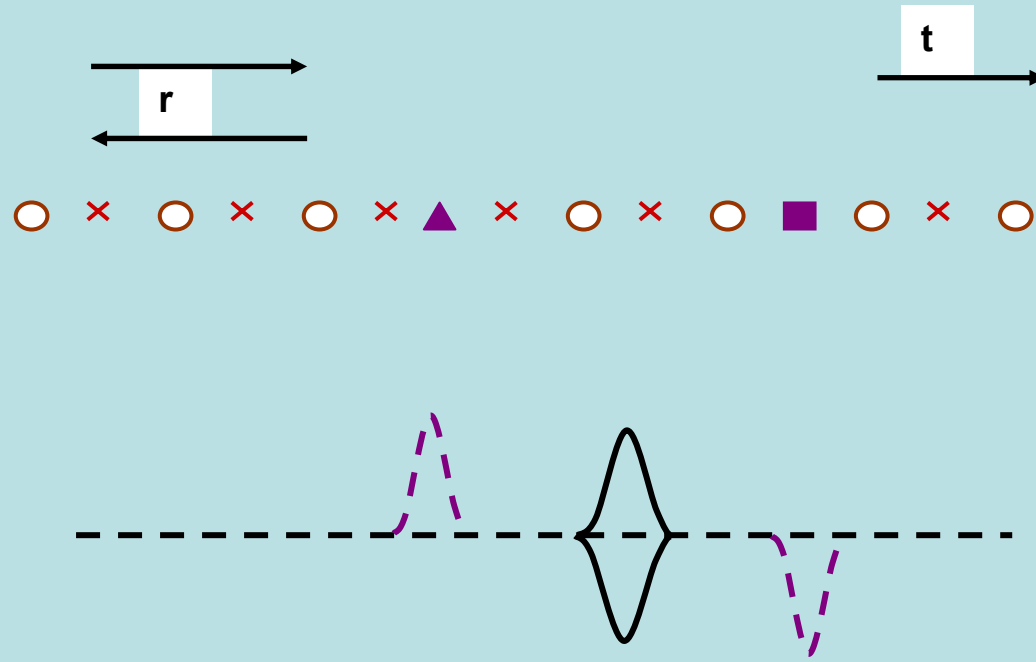
Dependence of effective interdefect distance on wave-vector



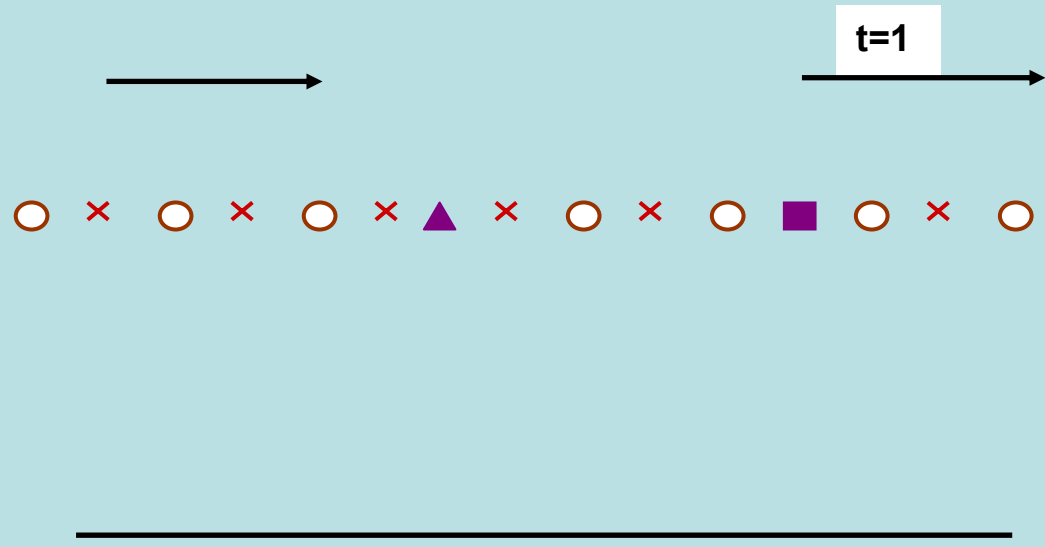
Dependence of reflection coefficient on wave-vector



Hidden Defect Pair



In continuum limit – homogeneous structure



Other examples of “dark matter” objects in systems without inversion center:

- **HDP in the models with continuum potentials (generalized Kronig-Penney model, pseudopotentials)**
- **Quantum well with defects located at the heterointerfaces**

There's is Plenty of Room at the Bottom.

R. Feinman, 1959

Nanotechnology – a lot of applications and economics, science and new knowledge, and some room for...game!